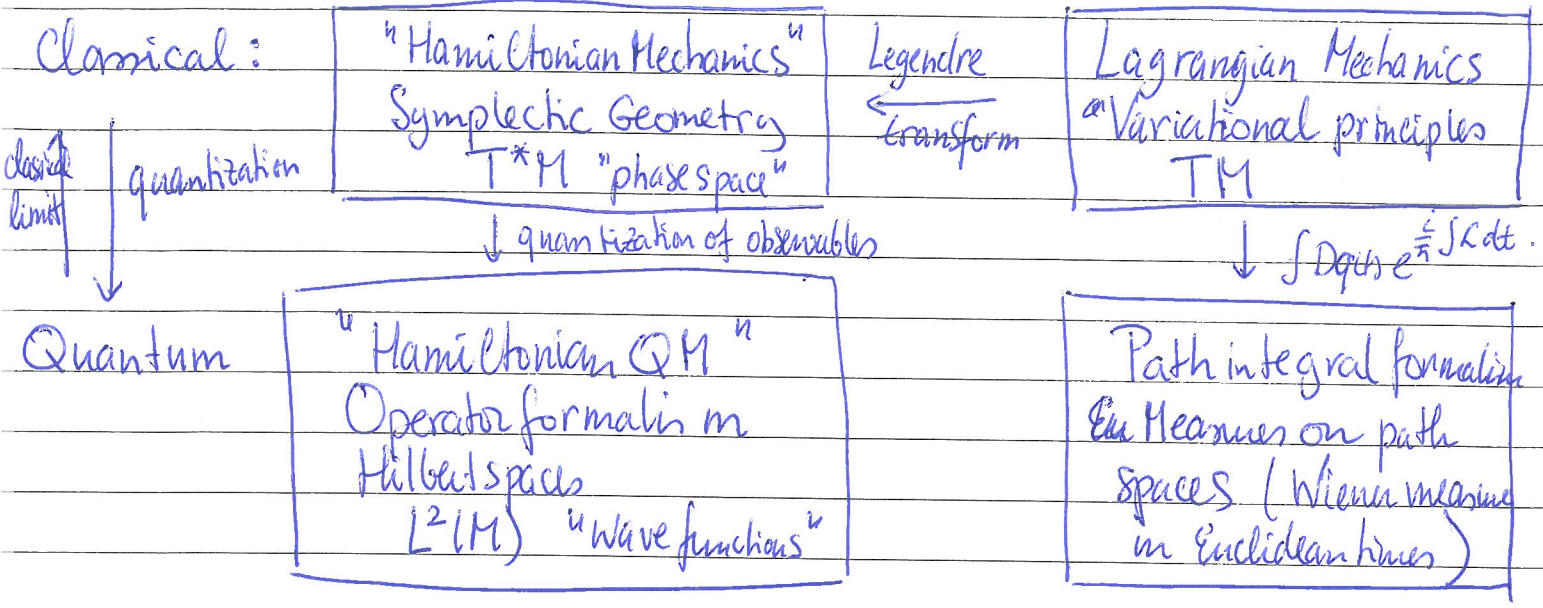


Before we do QFT, let's focus on finite dimensional systems (Quantum mechanics) The classical theory (called "classical" or "Newtonian" mechanics) can be formulated in 2-ways:



We will start in the upper right corner. Let's do the example of mechanics on a general manifold M . We shall (at some point) need a metric g on M , so let's focus on a Riemannian manifold (M, g) .

Def: A "Lagrangian function" is a C^1 -function $L: TM \rightarrow \mathbb{R}$.

Such a Lagrangian function leads to an action functional:

$$S'(y) : \int_{t_0}^{t_1} (y^* L) dt = \int_{t_0}^{t_1} L(y(t), \dot{y}(t)) dt$$

for $y: I \rightarrow M$, $I = [t_0, t_1]$ (we allow $t_0 = -\infty, t_1 = \infty \dots$)
The variational problem for S' now is: find the extrema of S' .

Recall for a finite dimensional manifold X , the extrema of a C^1 -function $f: X \rightarrow \mathbb{R}$, are the critical points $x_0 \in X$ where $df(x_0) = 0$. In other words

$$\frac{d}{d\varepsilon} f(x_0 + \varepsilon X) = 0 \quad \forall X \in T_{x_0} X.$$

The problem here now is the fact that we have to view S' as a function

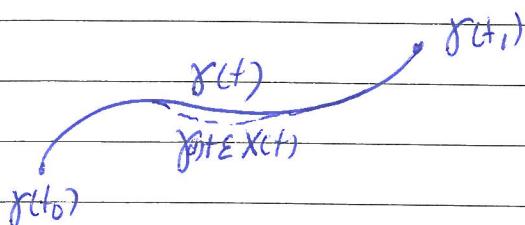
$$S': C^2(I; M) \longrightarrow \mathbb{R},$$

and the space $C^2(I; M)$ of C^2 -paths is - at least - an ∞ -D manifold! In the notes wI have worked hard to make $C^2(I; M)$ into a Banach manifold. Then we can require $dS' = 0$ to find the extrema... Here we take a different route.

Def: A path $\gamma(t)$ is an extremum of S' when

$$\left. \frac{d}{d\varepsilon} S'(\gamma + \varepsilon X) \right|_{\varepsilon=0} = 0$$

for any compact perturbation of γ .



compact = compact support.

Question: what is a compact perturbation precisely?

Answer: if we really construct $C^2(I; M)$ as a Banach manifold, a compact perturbation would be an element of $\Gamma_c^\infty(I; \gamma^* TM)_0$ = "compactly supported vector fields along $\gamma(t)$ ".

For $X = \mathbb{R}^n$, this is easy; in this case $C^2(I; \mathbb{R}^n)$ is a Banach space with norm

$$\|\gamma\|_{C^2} := \sup_{x \in I} \|\gamma(t)\| + \sup_{x \in I} \|\gamma'(t)\| + \sup_{x \in I} \|\gamma''(t)\|.$$

and we can use the linear structure: a compact perturbation is just a $X \in C_c^2(I; \mathbb{R}^n)$.

Def: $\gamma: I \rightarrow M$ is a critical point for S' if $\forall C^1$ -families of paths $\gamma_s: [-\epsilon, \epsilon] \times I \rightarrow M$ with $\gamma(0) = \gamma$ we have

$$\frac{d}{ds} S(\gamma_s) = 0 \quad \text{and } \gamma_s(t_0) = \gamma(t_0) \\ \gamma_s(t_1) = \gamma(t_1)$$

Prop: For $X = \mathbb{R}^n$, the extrema of the action are exactly the solutions of the Euler-Lagrange equations:

$$\frac{\partial L}{\partial x^i} - \frac{d}{dt} \frac{\partial L}{\partial \dot{x}^i} = 0 \quad \text{where } \gamma(t) = (x^1(t), \dots, x^n(t))$$

Pf: Let $\gamma_s(t) = (x_s^1(t), \dots, x_s^n(t))$ be a smooth variation of γ .

Define $X^i(t) = \frac{d}{ds} \Big|_{s=0} x_s^i(t)$. Then we have

$$\begin{aligned} \frac{d}{ds} \Big|_{s=0} S(\gamma_s(t)) &= \int_{t_0}^{t_1} \left(\frac{\partial L}{\partial x^i} X^i + \frac{\partial L}{\partial \dot{x}^i} \dot{X}^i \right) dt \\ &= \int_{t_0}^{t_1} \left(\frac{\partial L}{\partial x^i} X^i - \frac{d}{dt} \frac{\partial L}{\partial \dot{x}^i} X^i \right) dt + \underbrace{\left[\frac{\partial L}{\partial \dot{x}^i} X^i \right]_{t_0}^{t_1}}_{=0} \end{aligned}$$

We have: Lemma: if a continuous function $f(t)$ $t_0 \leq t \leq t_1$, satisfies $\int_{t_0}^{t_1} f(t)g(t) dt = 0$ for any g with $g(t_0) = g(t_1) = 0$ then $f \equiv 0$

□

For general M , we have the following

Lemma: If $\gamma: [t_0, t_1] \rightarrow M$ is an extremum of S' , for any sub interval $[t'_0, t'_1] \subset [t_0, t_1]$, $\gamma|_{[t'_0, t'_1]}$ is an extremum for S' .

Proof: let γ'_s be any ^{cpt} variation of $\gamma|_{[t'_0, t'_1]}$. Then, extending γ'_s by zero, this is also a variation of γ . Therefore

$$\frac{d}{ds} \Big|_{s=0} S'(\gamma'_s) = 0 \quad \text{as required}$$

□

Prop: For a general manifold, a path $\gamma: I \rightarrow M$ is an extremum of S' if it satisfies the Euler-Lagrange equations in local charts on M .

Prf: by the lemma, we can divide the interval I into smaller intervals ~~contain~~ whose image under γ is contained in a local chart. \square

Rk: This is only if and only if, if $\gamma(I)$ lies in 1 coordinate chart! (this because the lemma gives γ critical $\Rightarrow \gamma|_I$ critical.) so only necessary conditions.

Also, solving E-L in one coordinate system might be harder than solving it in another... Hamiltonian formalism!

Examples: (i) (zero order Lagrangians) suppose $L: TM \rightarrow \mathbb{R}$ is the pull-back of a function f on M .

$$(E-L) : \frac{\partial L}{\partial x^i}(\gamma(t)) = 0.$$

so γ should lie on the critical set of L on M .
 generic $L \Rightarrow$ critical pb are isolated \Rightarrow uninteresting

(ii) (first order Lagrangians) Suppose L is linear on the fibers of TM : $L(x, \xi) = L_0(x) + \alpha_x(\xi)$, $\alpha \in \Omega^1(M)$

$$(E-L) : \frac{\partial L_0}{\partial x^i} + \frac{\partial \alpha}{\partial x^i} - \frac{d}{dt} \alpha^i = 0$$

$$\frac{\partial L_0}{\partial x^i}(\gamma(t)) + \frac{\partial \alpha}{\partial x^i}(\gamma(t), \dot{\gamma}(t)) - \frac{d}{dt} \alpha^i(\gamma(t), \dot{\gamma}(t)) \quad \alpha = \alpha_i^j \xi^j$$

$$= \frac{\partial L_0}{\partial x^i}(\gamma(t)) + \left(\frac{\partial \alpha_j^i}{\partial x^k}(\gamma(t)) \dot{\gamma}^k - \frac{\partial \alpha^i}{\partial x^j} \frac{d\dot{\gamma}^j}{dt} \right)$$

coordinate free

$$= \frac{dL_0}{dt} + i_{\dot{\gamma}} d\alpha = 0$$

Conditions for a minimum

-6-

In finite dimensions, $df(x_0) = 0$ can mean that we have a ^(local) minimum, maximum or a saddle point. To distinguish these, we compute the Hessian $\frac{\partial^2 f}{\partial x^i \partial x^j}(x_0) = H_{ij}$. The cases are distinguished

$$\text{by } \det H_{ij} = \begin{cases} < 0 & \text{max} \\ > 0 & \text{min} \\ = 0 & \text{saddle} \end{cases}$$

Therefore, to investigate the behaviour of our action, we consider its second variation:

$$\frac{d^2 S}{ds^2}(\gamma_s) = \sum_{i,j=1}^n \int_{t_0}^{t_1} \left\{ \frac{\partial^2 L}{\partial x^i \partial x^j} \dot{x}^i \dot{x}^j + 2 \frac{\partial^2 L}{\partial x^i \partial \dot{x}^j} \dot{x}^i \ddot{x}^j + \frac{\partial^2 L}{\partial \dot{x}^i \partial \dot{x}^j} \ddot{x}^i \ddot{x}^j \right\} dt$$
$$\stackrel{\text{Integration by parts}}{=} \sum_{i,j} \int_{t_0}^{t_1} \left(Q_{ij} \dot{x}^i \dot{x}^j + P_{ij} \ddot{x}^i \ddot{x}^j \right) dt \quad (*)$$

$$\text{with } Q_{ij} = \frac{\partial^2 L}{\partial \dot{x}^i \partial \dot{x}^j}, \quad P_{ij} = \frac{\partial^2 L}{\partial \dot{x}^i \partial \ddot{x}^j} - \frac{d}{dt} \frac{\partial^2 L}{\partial x^i \partial \ddot{x}^j}$$

$$\text{where we assume } \frac{\partial^2 L}{\partial x^i \partial \ddot{x}^j} = \frac{\partial^2 L}{\partial \ddot{x}^j \partial x^i}$$

Claim: the dominant term in (*) is given by Q.

this is because we can view (*) itself as a quadratic functional on the space of paths $X(t)$, and call this I_X . Clearly $I_X(0) = 0$ and we are interested in the variation of I_X for small X . Now we have

\dot{X} small $\Rightarrow X$ small, but not the other way around, we can easily construct a small X (in the C^0 -norm) with large \dot{X} . Therefore we find Legendre's condition

$$\det \left(\frac{\partial^2 L}{\partial \dot{x}^i \partial \dot{x}^j} \right) \geq 0 \quad \text{along } \gamma(t)$$

Consider the variational problem associated to \mathbb{I}_γ : its Euler-Lagrange equations are given by

$$\sum_{i=1}^n \left\{ \frac{\partial^2 L}{\partial x^i \partial x^i} \dot{x}^i + \frac{\partial^2 L}{\partial x^i \partial \xi^i} \dot{x}^i - \frac{d}{dt} \left(\frac{\partial^2 L}{\partial x^i \partial \xi^i} \dot{x}^i + \frac{\partial^2 L}{\partial \xi^i \partial \xi^i} \dot{x}^i \right) \right\} = 0$$

Jacobi equations

Prop if γ_s is a smooth family of solutions of E-L for \mathcal{S} ,
 $X(t) := \frac{d}{ds} \Big|_{s=0} \gamma_s(t)$ is a solution of the Jacobi-equations

Proof differentiate E-L. □

Thm: \mathbb{I}_γ is positive definite if

- (i) $\det \left(\frac{\partial^2 L}{\partial \xi^i \partial \xi^i} \right) > 0$ along $\gamma(t)$
- (ii) the interval $(t_0, t_1]$ does not contain a point conjugate to t_0 .

Rk: Legendre tried to prove the theorem without (ii).

Examples (continued) (iii) quadratic Lagrangians

Because of Legendre's condition, we are forced to take

$$\begin{aligned} L(x, \xi) &= \frac{1}{2} g_x(\xi, \xi), \quad g \text{ a Riemannian metric.} \\ &= \frac{1}{2} \|\xi\|^2 \implies \mathcal{S}'(\gamma) = \frac{1}{2} \int_{t_0}^{t_1} \|\dot{\gamma}(t)\|^2 dt \end{aligned}$$

Prop: The Euler-Lagrange equations for L are given by the geodesic equations $\nabla_{\dot{\gamma}} \dot{\gamma} = 0$.

Proof: let ∇ be the Levi-Civita connection associated to g , determined by the equation

$$d g(X, Y) = g(\nabla X, Y) + g(X, \nabla Y)$$

Let γ_s be a compact perturbation of γ_0 . Then

$$\begin{aligned}
 \frac{d}{ds} S(\gamma_s(t)) &= \frac{1}{2} \frac{d}{ds} \int_{t_0}^{t_1} g\left(\frac{d\gamma_s}{dt}, \frac{d\gamma_s}{dt}\right) dt \\
 &= \int_{t_0}^{t_1} g\left(\nabla_{\frac{d\gamma_s}{ds}}\left(\frac{d\gamma}{dt}\right), \frac{d\gamma}{dt}\right) dt \\
 &= \int_{t_0}^{t_1} g\left(\nabla_{\frac{d\gamma}{dt}}\left(\frac{d\gamma}{ds}\right), \frac{d\gamma}{dt}\right) dt \\
 &= \int_{t_0}^{t_1} \left\{ \frac{d}{dt} g\left(\frac{d\gamma}{ds}, \frac{d\gamma}{dt}\right) - g\left(\frac{d\gamma}{ds}, \nabla_{\frac{d\gamma}{dt}} \frac{d\gamma}{dt}\right) \right\} dt \\
 &= - \int_{t_0}^{t_1} g\left(\frac{d\gamma}{ds}, \nabla_{\frac{d\gamma}{dt}} \frac{d\gamma}{dt}\right) dt \quad \square
 \end{aligned}$$

Rk: if we include a potential $V \in C^\infty(M)$, we get

$$\nabla_{\frac{d\gamma}{dt}} \frac{d\gamma}{dt} = -\text{grad}(V)$$