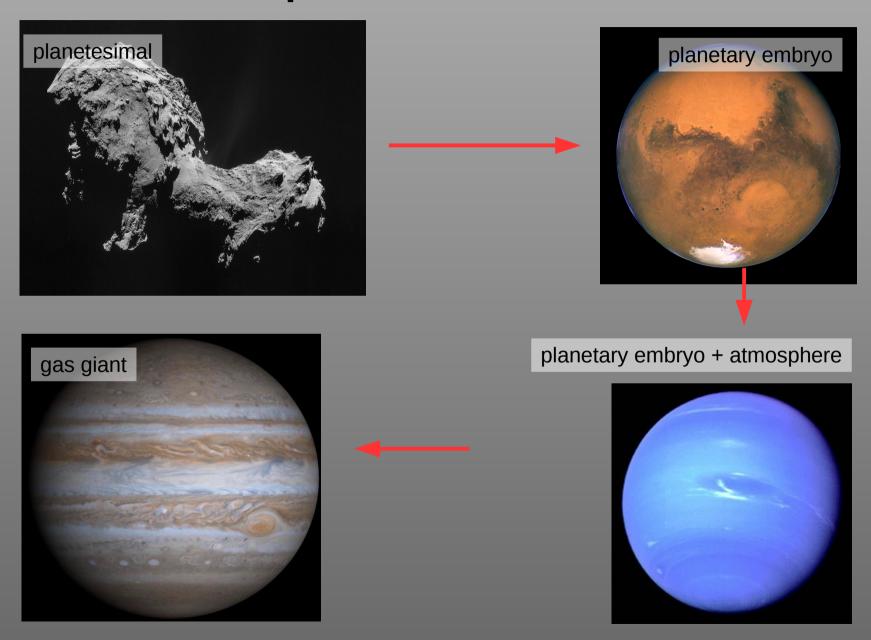
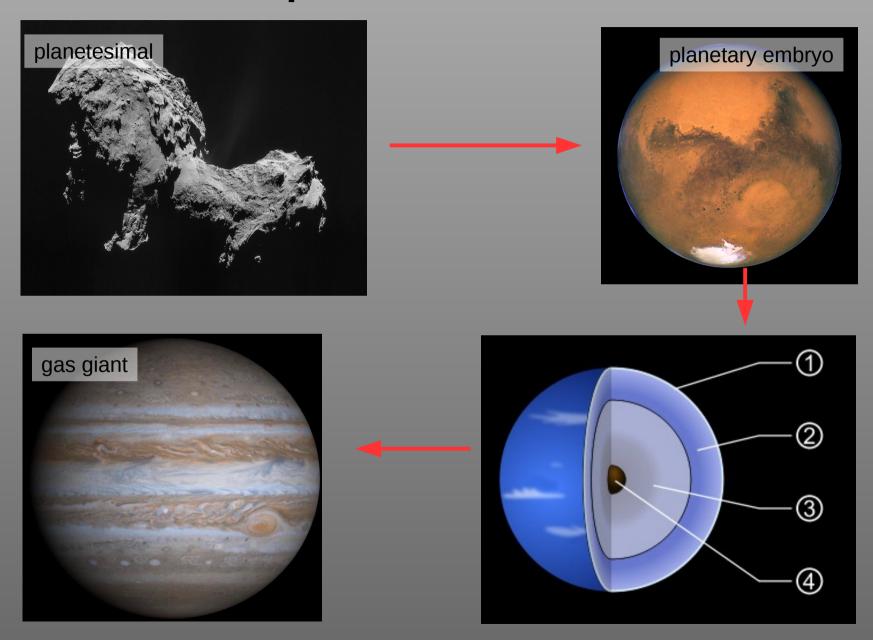
# L12: Protoplanet growth, planet atmospheres, & giant planet formation



# L12: Protoplanet growth, planet atmospheres, & giant planet formation



# Features oligarchic growth

#### **Towards oligarchy**

- diverging (runaway) growth w/i same zone
- converging (normal) growth for different zones

#### 2 components

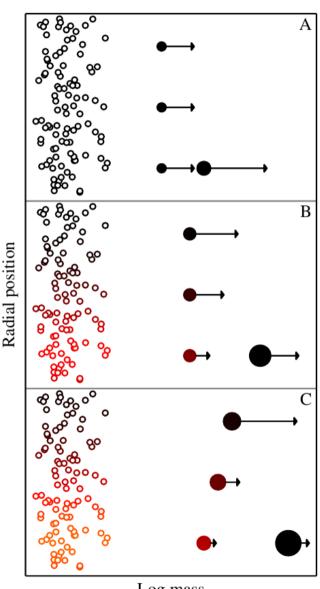
- planetesimals (dominate  $\Sigma$  initially)
- embryos (dominate dynamics)

#### **During oligarchy**

embryos feast on planetesimals, but also merge; feeding zone stays several  $R_{\text{Hill}}$ .

#### Slower than R.G.

but can still feature large  $\Theta$  especially when planetesimals are damped (by gas).



Log mass

# Velocity regimes

#### Dispersiondominated regime

Relative velocity (v<sub>∞</sub>) determined by eccentric motion of planetesimal

$$v_{\infty} = ev_{K}$$

# **Shear-dominated** regime

v<sub>∞</sub>determined by Keplerian shear

$$v_{\infty} = (3/2)b\Omega_{K}$$

#### **Headwind regime**

v<sub>w</sub>determined by sub-Keplerian headwind gas

$$v_{\infty} = \eta v_{K}$$

#### **Planetesimal Accretion**

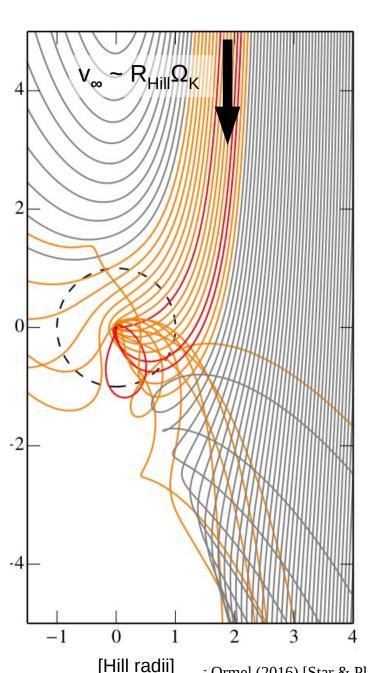
gas drag damps eccentricity on long timescales ( $\tau_p >> 1$ )

#### **Pebble Accretion**

gas drag acts during encounter ( $t_{stop}$  small)

# Shear-dominated interactions

w/o gas drag



#### Hill accretion (shear-dominated; planetesimals)

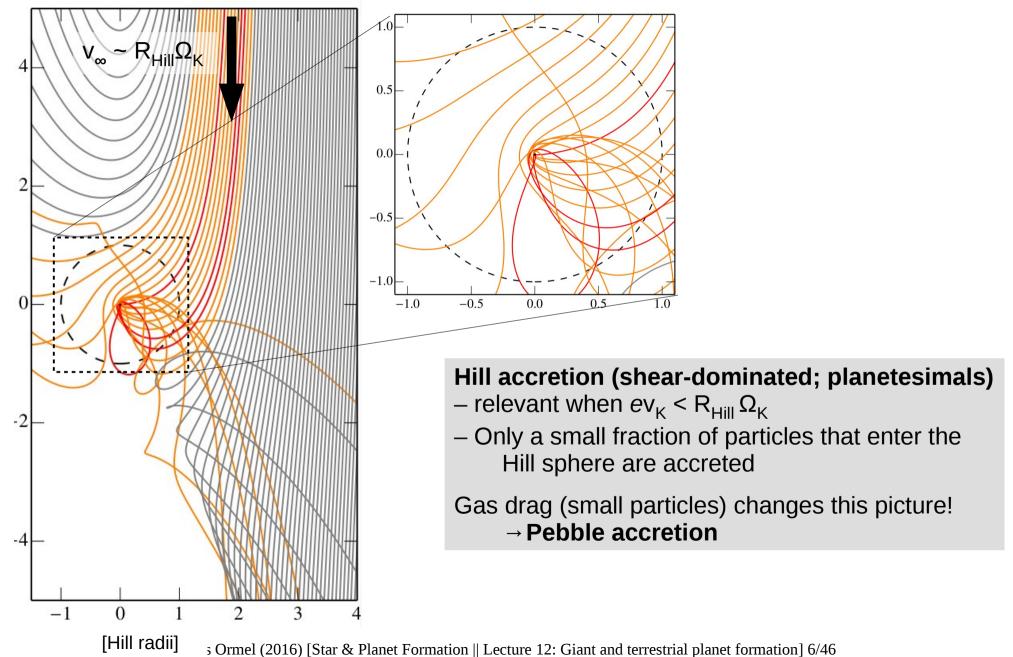
- relevant when  $ev_K < R_{Hill} \Omega_K$
- Only a small fraction of particles that enter the Hill sphere are accreted

Gas drag (small particles) changes this picture!

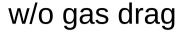
→ Pebble accretion

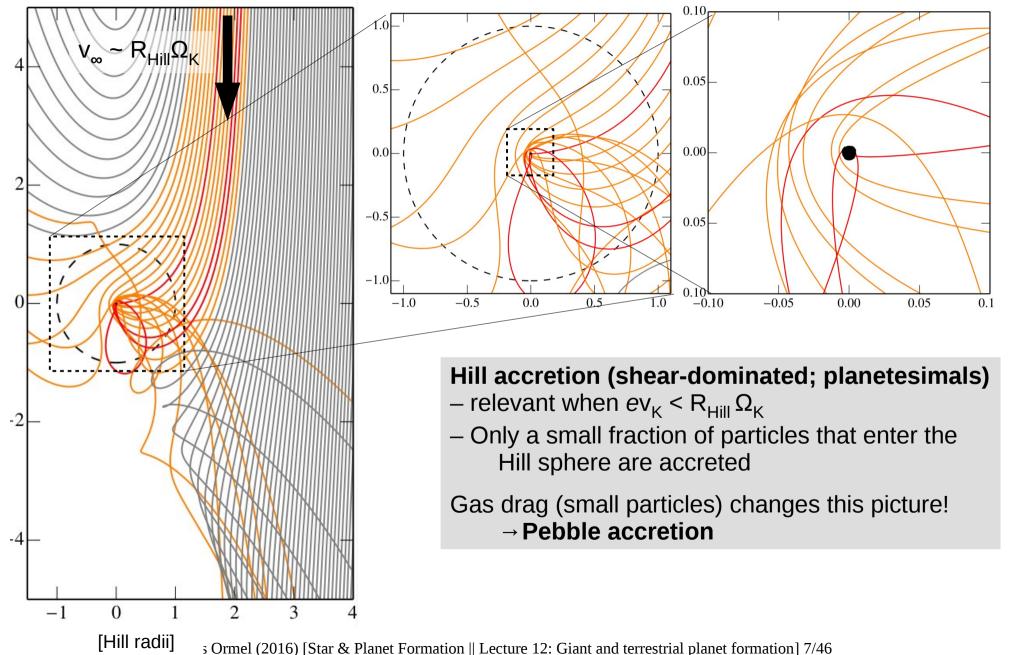
# Shear-dominated interactions

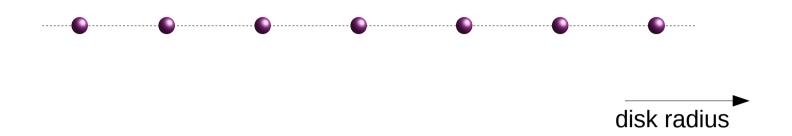
w/o gas drag

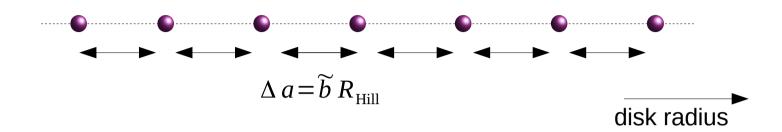


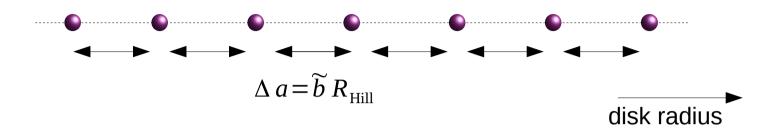
# Shear-dominated interactions





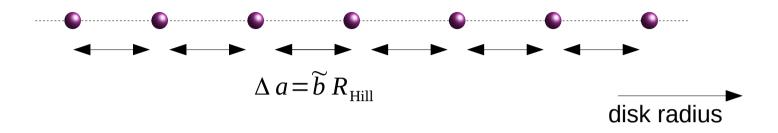






#### **Isolation mass**

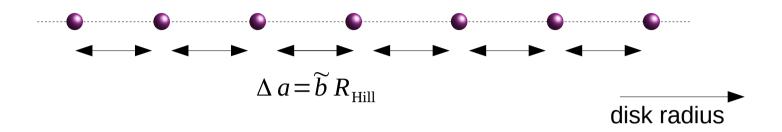
$$M_{\rm iso}(R_{\rm Hill}) = 2 \pi a \Sigma \Delta a$$



#### **Isolation mass**

$$M_{\rm iso}(R_{\rm Hill}) = 2 \pi a \Sigma \Delta a$$

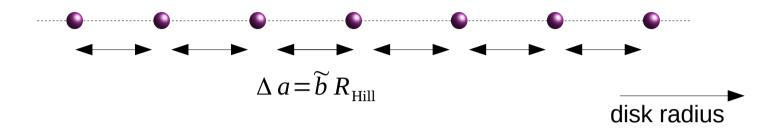
$$M_{\rm iso} = \frac{(2\pi \tilde{b}\Sigma r^2)^{3/2}}{(3M_{\oplus})^{1/2}} \simeq 0.25 \ M_{\oplus} \left(\frac{\tilde{b}}{10}\right)^{3/2} \left(\frac{\Sigma}{10}\right)^{-3/2} \left(\frac{a}{{
m AU}}\right)^3$$



#### **Isolation mass**

$$M_{\rm iso}(R_{\rm Hill}) = 2 \pi a \Sigma \Delta a$$

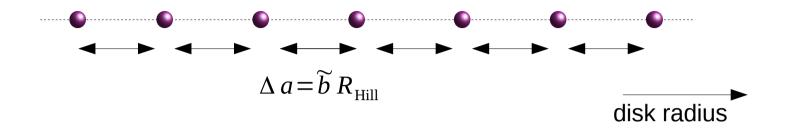
$$M_{\rm iso} = \frac{(2\pi \tilde{b} \Sigma \mathbf{a^2})^{3/2}}{(3M_{\oplus})^{1/2}} \simeq 0.25 \ M_{\oplus} \left(\frac{\tilde{b}}{10}\right)^{3/2} \left(\frac{\Sigma}{10}\right)^{-3/2} \left(\frac{a}{{
m AU}}\right)^3$$



#### **Isolation mass**

$$M_{\rm iso}(R_{\rm Hill}) = 2 \pi a \Sigma \Delta a$$

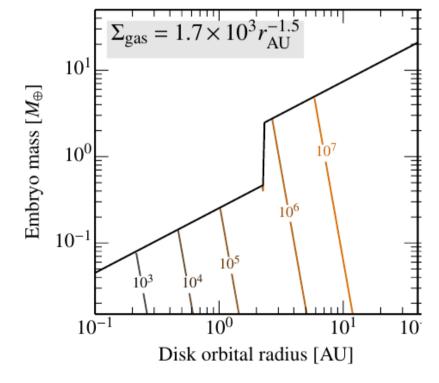
$$M_{\rm iso} = \frac{(2\pi \tilde{b} \Sigma \mathbf{a^2})^{3/2}}{(3M_{\oplus})^{1/2}} \simeq 0.25 \ M_{\oplus} \left(\frac{\tilde{b}}{10}\right)^{3/2} \left(\frac{\Sigma}{10}\right)^{+3/2} \left(\frac{a}{{
m AU}}\right)^3$$

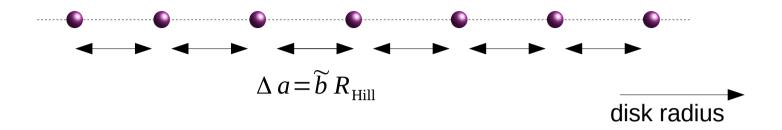


#### **Isolation mass**

$$M_{\rm iso}(R_{\rm Hill}) = 2 \pi a \Sigma \Delta a$$

$$M_{\rm iso} = \frac{(2\pi \tilde{b} \Sigma \mathbf{a^2})^{3/2}}{(3M_{\oplus})^{1/2}} \simeq 0.25 \ M_{\oplus} \left(\frac{\tilde{b}}{10}\right)^{3/2} \left(\frac{\Sigma}{10}\right)^{+3/2} \left(\frac{a}{{
m AU}}\right)^3$$

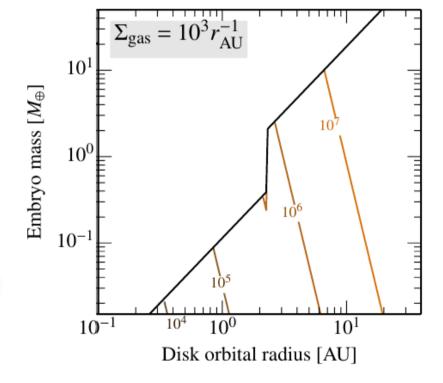




#### **Isolation mass**

$$M_{\rm iso}(R_{\rm Hill}) = 2 \pi a \Sigma \Delta a$$

$$M_{\rm iso} = \frac{(2\pi \tilde{b} \Sigma \mathbf{a^2})^{3/2}}{(3M_{\oplus})^{1/2}} \simeq 0.25 \ M_{\oplus} \left(\frac{\tilde{b}}{10}\right)^{3/2} \left(\frac{\Sigma}{10}\right)^{+3/2} \left(\frac{a}{{
m AU}}\right)^3$$



# Giant Planet formation

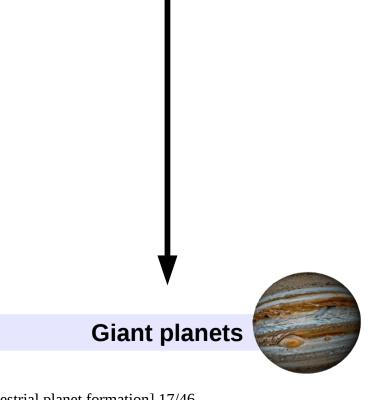
**Disk instability model** gravitational instab. gas

- Toomre-Q < 1
- efficient cooling

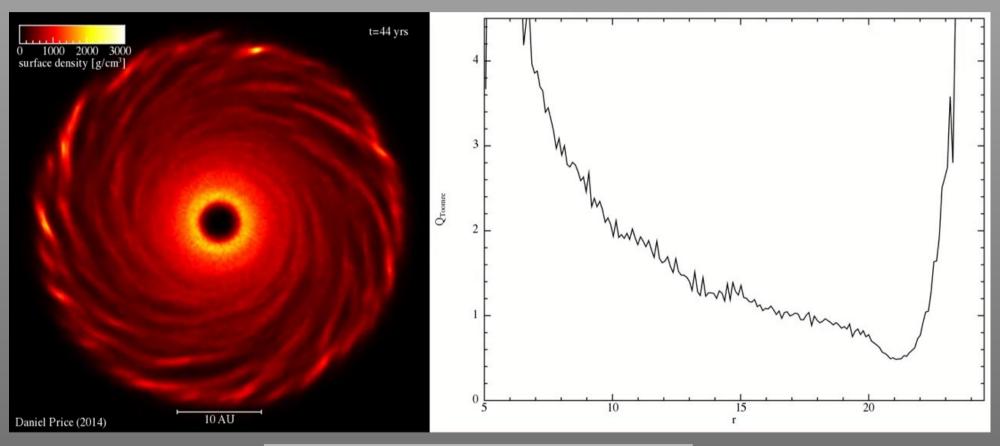
# Giant Planet formation

**Disk instability model** gravitational instab. gas

- Toomre-Q < 1
- efficient cooling

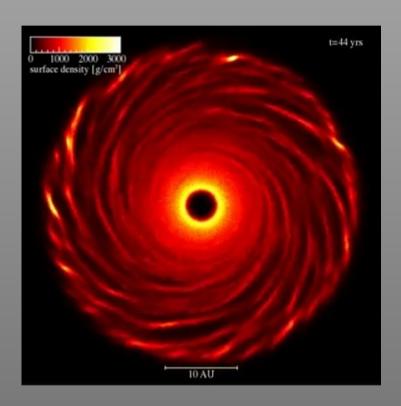


# Disk instability



Price (2014)
https://www.youtube.com/watch?v=hngA5CKIs58

# Efficient cooling

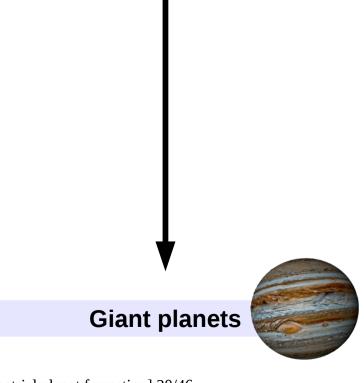


Price (2014)
https://www.youtube.com/watch?v=\_JgwlWDL3aw

#### Disk instability model

(Lecture 9)

- Toomre-Q < 1
- efficient cooling



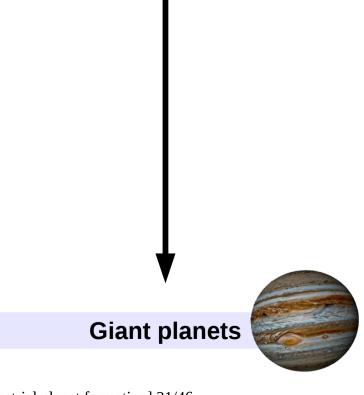
#### **Planetesimals**

- sticking (L8)
- GI instability (L9)

#### Disk instability model

(Lecture 9)

- Toomre-Q < 1</p>
- efficient cooling



#### **Planetesimals**

- sticking (L8)
- GI instability (L9)

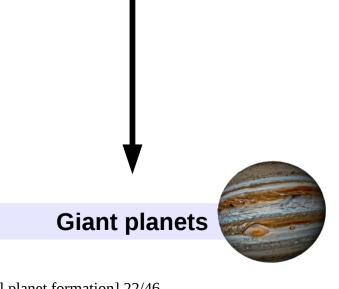
Runway and oligarchic growth, pebble accr. (L11)

Planetary embryos isolated

#### Disk instability model

(Lecture 9)

- Toomre-Q < 1</p>
- efficient cooling



#### **Planetesimals**

- sticking (L8)
- GI instability (L9)

**\** 

Runway and oligarchic growth, pebble accr. (L11)

# Planetary embryos isolated

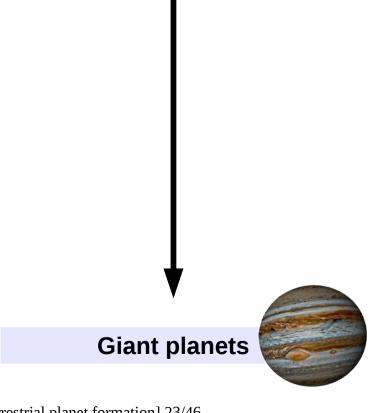
small embryos

Gas disk disappears (Lecture 12)

Terrestrial planet orbit crossing

# **Disk instability model** (Lecture 9)

- \_ Toomre-Q < 1
- efficient cooling



#### **Planetesimals**

- sticking (L8)
- GI instability (L9)

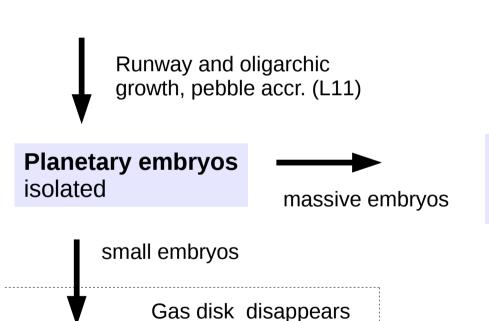
**Terrestrial planet** 

orbit crossing

Disk instability model

(Lecture 9)

- Toomre-Q < 1</p>
- efficient cooling



(Lecture 12)

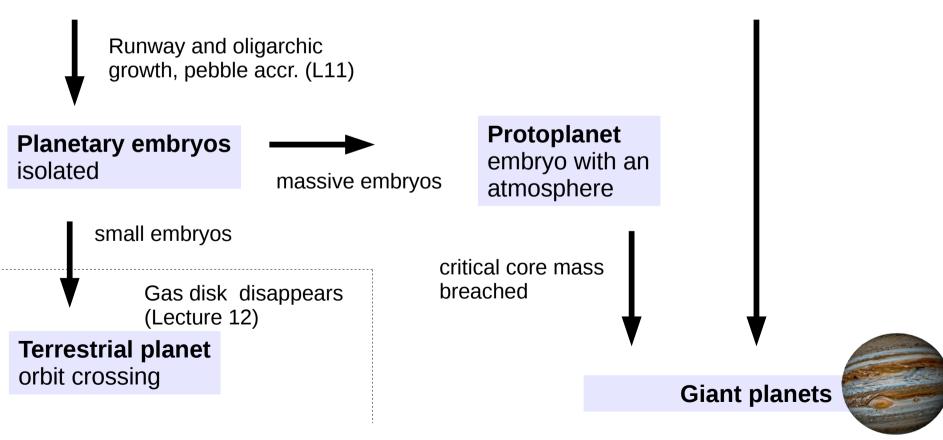
**Protoplanet** embryo with an atmosphere

Giant planets

# Planetesimals – sticking (L8) – GI instability (L9)

# **Disk instability model** (Lecture 9) – Toomre-Q < 1

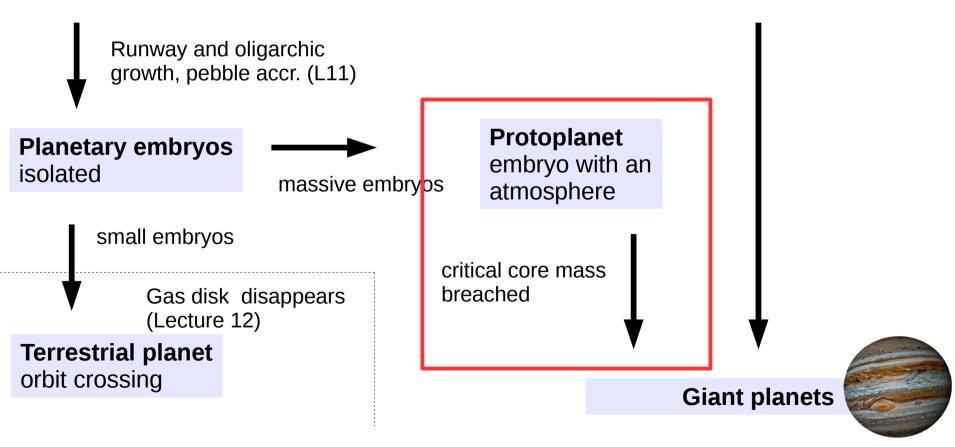
efficient cooling



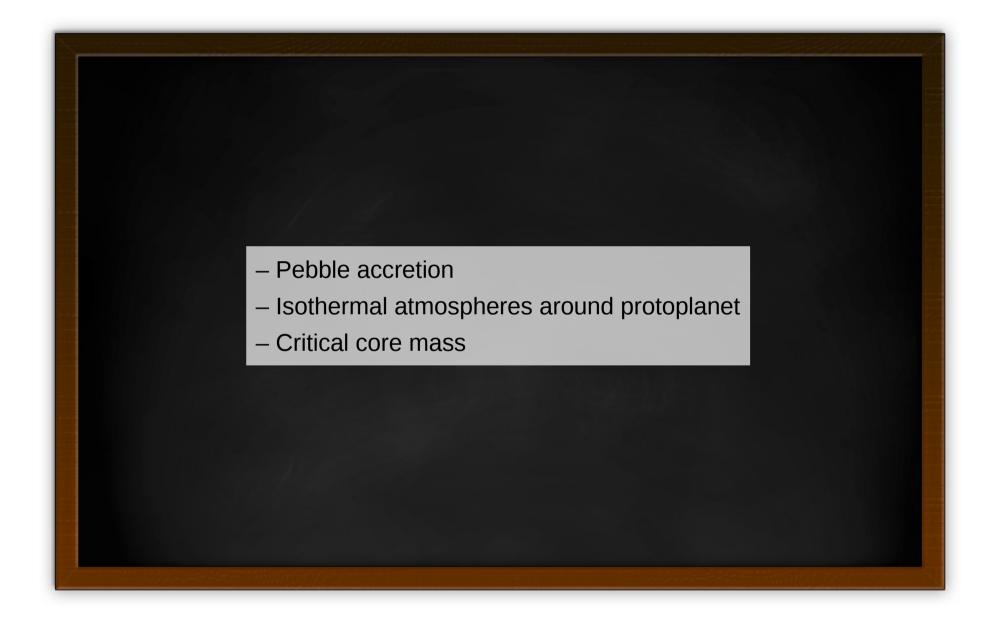
# Planetesimals – sticking (L8) – GI instability (L9)

# Disk instability model (Lecture 9)

- Toomre-Q < 1</p>
- efficient cooling



# Blackboard



#### Realistic atmosphere models

solve the stellar-structure equations:

#### **Energy transport**: $\nabla = \min(\nabla_{rad}, \nabla_{ad})$

- $-\nabla_{rad}$ :transport by radiation
- $-\nabla_{ad}$  :transport by convection

- boundary conditions
- Luminosity source

Continuity 
$$rac{\partial M_{< r}}{\partial r} = 4\pi G r^2 
ho$$

Hydrostatic balance 
$$\frac{\partial P}{\partial r} = -\rho \frac{GM_{< r}}{r^2}$$

Energy transport 
$$\frac{\partial T}{\partial r} = \frac{\partial P}{\partial r} \frac{T}{P} \nabla$$

Luminosity conservation 
$$\frac{\partial L}{\partial r} = 4\pi r^2 \rho \left( \epsilon - T \frac{dS}{dt} \right)$$

E.O.S. 
$$P = P(T, \rho)$$

$$\nabla_{\text{rad}} = \frac{3 \,\kappa \,L \,P}{64 \,\pi \,\sigma_{\text{sb}} \,G M_{\,<\text{r}} \,T^4} \qquad \nabla_{\text{ad}} = \left(\frac{\text{d} \log T}{\text{d} \log P}\right)_{\text{ad}}$$

#### Realistic atmosphere models

solve the stellar-structure equations:

#### **Energy transport**: $\nabla = \min(\nabla_{rad}, \nabla_{ad})$

- $-\nabla_{rad}$ :transport by radiation
- $-\nabla_{ad}$  :transport by convection

- boundary conditions
- Luminosity source

Continuity 
$$rac{\partial M_{< r}}{\partial r} = 4\pi G r^2 
ho$$

Hydrostatic balance 
$$\frac{\partial P}{\partial r} = -\rho \frac{GM_{< r}}{r^2}$$

Energy transport 
$$\frac{\partial T}{\partial r} = \frac{\partial P}{\partial r} \frac{T}{P} \nabla$$

Luminosity conservation 
$$\frac{\partial L}{\partial r} = 4\pi r^2 \rho \left( \epsilon - T \frac{dS}{dt} \right)$$

E.O.S. 
$$P = P(T, \rho)$$

$$\nabla_{\text{rad}} = \frac{3 \,\kappa \,L \,P}{64 \,\pi \,\sigma_{\text{sb}} \,G M_{\,<\text{r}} \,T^4} \qquad \nabla_{\text{ad}} = \left(\frac{\text{d} \log T}{\text{d} \log P}\right)_{\text{ad}}$$

#### Realistic atmosphere models

solve the stellar-structure equations:

#### **Energy transport**: $\nabla = \min(\nabla_{rad}, \nabla_{ad})$

- $-\nabla_{rad}$ :transport by radiation
- $-\nabla_{ad}$  :transport by convection

- boundary conditions
- Luminosity source

Continuity 
$$rac{\partial M_{< r}}{\partial r} = 4\pi G r^2 
ho$$

Hydrostatic balance 
$$\frac{\partial P}{\partial r} = -\rho \frac{GM_{< r}}{r^2}$$

Energy transport 
$$\frac{\partial T}{\partial r} = \frac{\partial P}{\partial r} \frac{T}{P} \nabla$$

Luminosity conservation 
$$\frac{\partial L}{\partial r} = 4\pi r^2 \rho \left( \epsilon - T \frac{dS}{dt} \right)$$

E.O.S. 
$$P = P(T, \rho)$$

$$\nabla_{\text{rad}} = \frac{3 \,\kappa \,L \,P}{64 \,\pi \,\sigma_{\text{sb}} \,GM_{\,<\text{r}} \,T^4} \qquad \nabla_{\text{ad}} = \left(\frac{\text{d} \log T}{\text{d} \log P}\right)_{\text{ad}}$$

#### Realistic atmosphere models

solve the stellar-structure equations:

#### **Energy transport**: $\nabla = \min(\nabla_{rad}, \nabla_{ad})$

- $-\nabla_{rad}$ :transport by radiation
- $-\nabla_{ad}$  :transport by convection

- boundary conditions
- Luminosity source

Continuity 
$$rac{\partial M_{< r}}{\partial r} = 4\pi G r^2 
ho$$

Hydrostatic balance 
$$\frac{\partial P}{\partial r} = -\rho \frac{GM_{< r}}{r^2}$$

Energy transport 
$$\frac{\partial T}{\partial r} = \frac{\partial P}{\partial r} \frac{T}{P} \nabla$$

Luminosity conservation 
$$\frac{\partial L}{\partial r} = 4\pi r^2 \rho \left( \epsilon - T \frac{dS}{dt} \right)$$

E.O.S. 
$$P = P(T, \rho)$$

$$\nabla_{\text{rad}} = \frac{3\kappa L P}{64 \pi \sigma_{\text{sb}} G M_{\text{cr}} T^4} \qquad \nabla_{\text{ad}} = \left(\frac{\text{d} \log T}{\text{d} \log P}\right)_{\text{ad}}$$

#### Realistic atmosphere models

solve the stellar-structure equations:

#### **Energy transport**: $\nabla = \min(\nabla_{rad}, \nabla_{ad})$

- $-\nabla_{rad}$ :transport by radiation
- $-\nabla_{ad}$  :transport by convection

- boundary conditions
- Luminosity source

Continuity 
$$rac{\partial M_{< r}}{\partial r} = 4\pi G r^2 
ho$$

Hydrostatic balance 
$$\frac{\partial P}{\partial r} = -\rho \frac{GM_{< r}}{r^2}$$

Energy transport 
$$\frac{\partial T}{\partial r} = \frac{\partial P}{\partial r} \frac{T}{P} \nabla$$

Luminosity conservation 
$$\frac{\partial L}{\partial r} = 4\pi r^2 \rho \left( \varepsilon - T \frac{dS}{dt} \right)$$

E.O.S. 
$$P = P(T, \rho)$$

$$\nabla_{\text{rad}} = \frac{3\kappa L P}{64 \pi \sigma_{\text{sb}} G M_{\text{cr}} T^4} \qquad \nabla_{\text{ad}} = \left(\frac{\text{d} \log T}{\text{d} \log P}\right)_{\text{ad}}$$

#### Realistic atmosphere models

solve the stellar-structure equations:

#### **Energy transport**: $\nabla = \min(\nabla_{rad}, \nabla_{ad})$

- $-\nabla_{rad}$ :transport by radiation
- $-\nabla_{ad}$ :transport by convection

- boundary conditions
- Luminosity source

Continuity 
$$\frac{\partial M_{< r}}{\partial r} = 4\pi G r^2 
ho$$

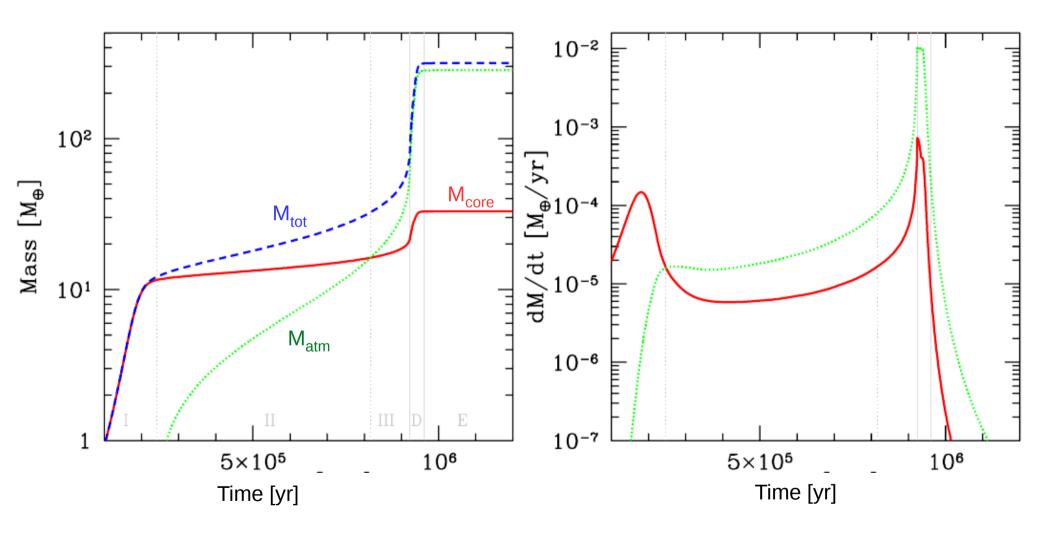
Hydrostatic balance 
$$\frac{\partial P}{\partial r} = -\rho \frac{GM_{< r}}{r^2}$$

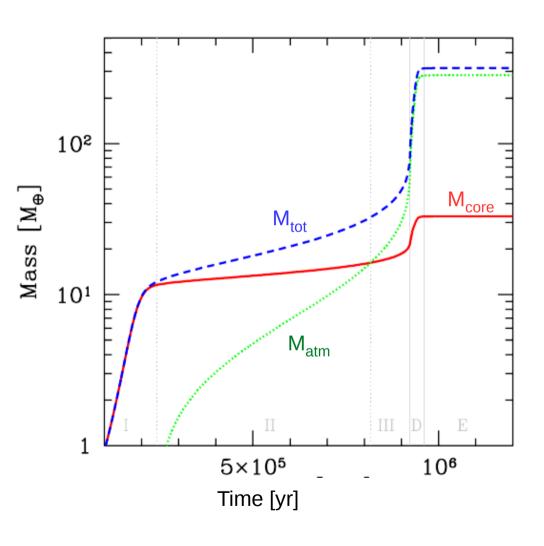
Energy transport 
$$\frac{\partial T}{\partial r} = \frac{\partial P}{\partial r} \frac{T}{P} \nabla$$

Luminosity conservation 
$$\frac{\partial L}{\partial r} = 4\pi r^2 \rho \left(\varepsilon\right) \left(T\frac{dS}{dt}\right)$$

E.O.S. 
$$P = P(T, \rho)$$

$$\nabla_{\text{rad}} = \frac{3\kappa L P}{64 \pi \sigma_{\text{sb}} G M_{\text{cr}} T^4} \qquad \nabla_{\text{ad}} = \left(\frac{\text{d} \log T}{\text{d} \log P}\right)_{\text{ad}}$$





#### Phase I

M<sub>tot</sub> increases steeply

#### **Phase II**

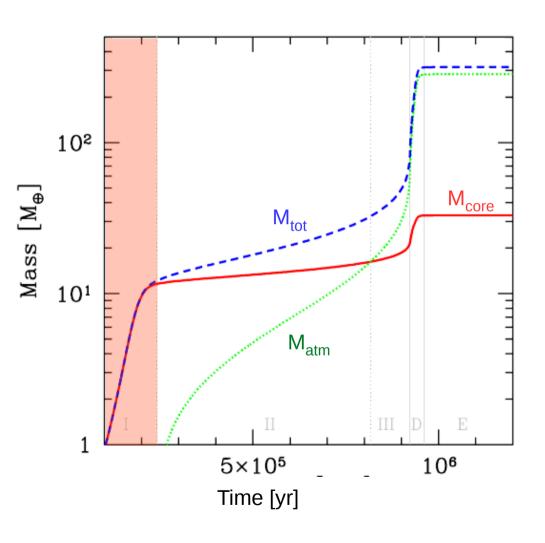
M<sub>tot</sub> slowly increases

#### **Phase III**

M<sub>tot</sub> rapidly increases

#### **Phase IV**

growth M<sub>tot</sub> stops



#### Phase I

M<sub>tot</sub> increases steeply

#### **Phase II**

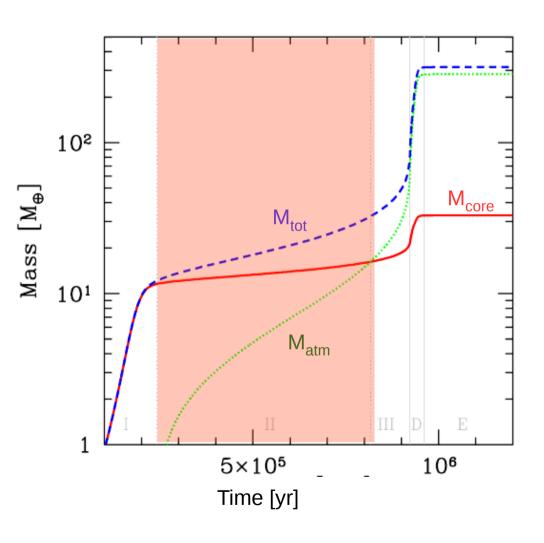
M<sub>tot</sub> slowly increases

#### **Phase III**

M<sub>tot</sub> rapidly increases

#### **Phase IV**

growth M<sub>tot</sub> stops



#### Phase I

 $M_{tot}$  increases steeply

#### **Phase II**

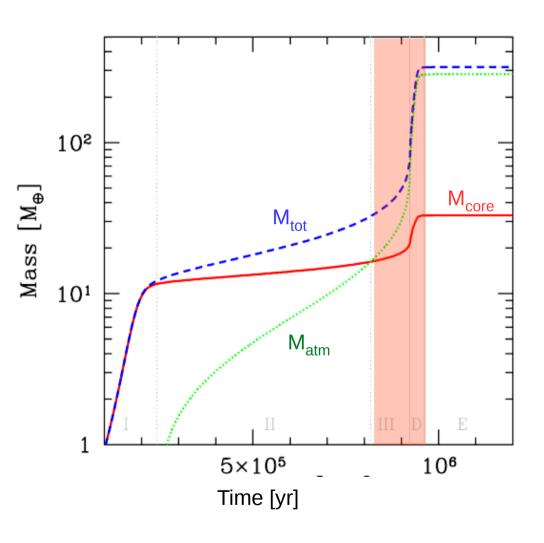
M<sub>tot</sub> slowly increases

#### **Phase III**

M<sub>tot</sub> rapidly increases

#### **Phase IV**

growth  $M_{tot}$  stops



#### Phase I

M<sub>tot</sub> increases steeply

#### **Phase II**

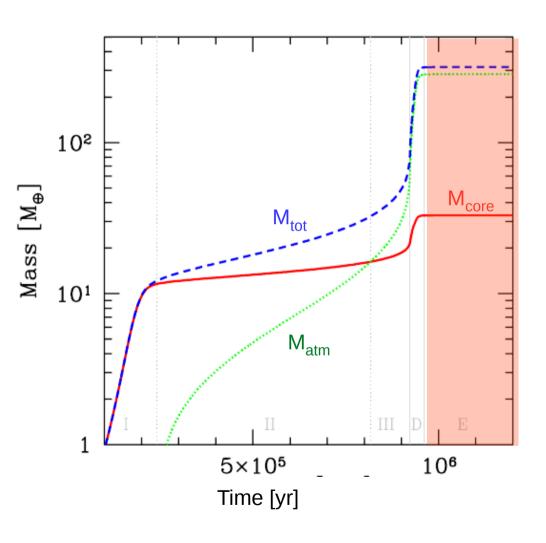
M<sub>tot</sub> slowly increases

#### **Phase III**

M<sub>tot</sub> rapidly increases

### Phase IV

growth M<sub>tot</sub> stops



#### Phase I

M<sub>tot</sub> increases steeply

#### **Phase II**

M<sub>tot</sub> slowly increases

#### **Phase III**

M<sub>tot</sub> rapidly increases

#### **Phase IV**

growth M<sub>tot</sub> stops

# Exercise 1.20

#### **Exercise 1.20 Radiative zero solution:** Assume the following:

- $\kappa$  and L are constant and define  $W \equiv 3\kappa L/64\pi\sigma_{\rm sb}$  (also constant);
- The gravitational mass interior to r,  $GM_{< r}$  can be approximated by the total mass of the planet+atmosphere,  $GM_{\text{tot}}$ , which is a constant;
- The atmosphere is radiatively supported:  $\nabla = \nabla_{\rm rad} = WP/GM_rT^4$ ;
- An ideal EOS,  $P = k_B \rho T / \mu$ .
- (a) Under these assumptions, show that Equation (1.58c) gives:

$$P = \frac{GM_{\text{tot}}T^4}{4W}; \qquad \rho = \frac{GM_{\text{tot}}\mu T^3}{4k_BW}$$
 (1.60)

where we neglected the boundary condition (the solutions are valid only in the "deep" atmosphere).

**(b)** Continue, by invoking Equations (1.58b) and (1.58c), to derive the atmosphere temperature and density profiles:

$$T(r) \simeq \frac{GM\mu}{4k_B} \frac{1}{r}; \qquad \rho(r) \simeq \frac{1}{W} \left(\frac{GM\mu}{4k_B}\right)^4 \frac{1}{r^3}.$$
 (1.61)

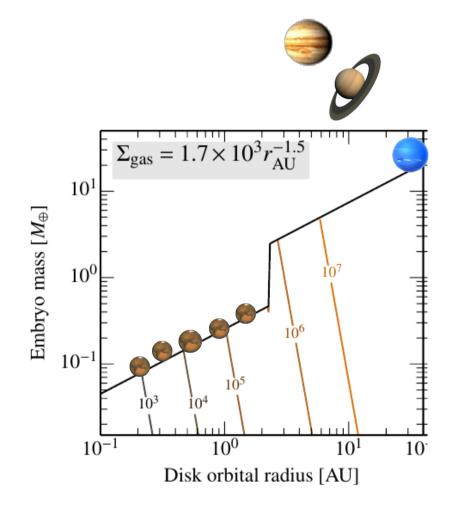
Integrating these gives the mass of the atmosphere:

$$M_{\text{atm}} = \frac{4\pi}{W} \left( \frac{GM_{\text{tot}}\mu}{4k_B} \right)^4 \Lambda, \tag{1.62}$$

# Assessment: oligarchic growth model

Key advantage oligarchic growth model/planetesimal accretion
Tailored to solar system

Drawbacks...



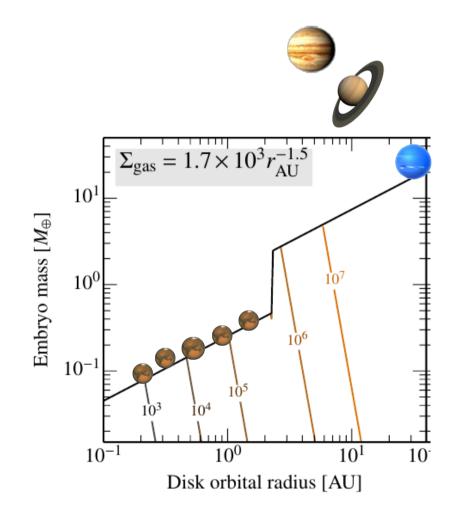
# Assessment: oligarchic growth model

# Key advantage oligarchic growth model/planetesimal accretion

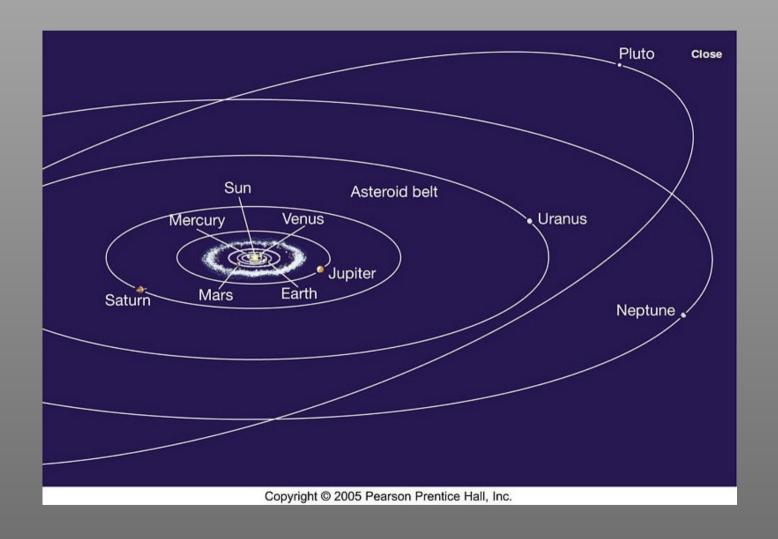
Tailored to solar system

#### Drawbacks...

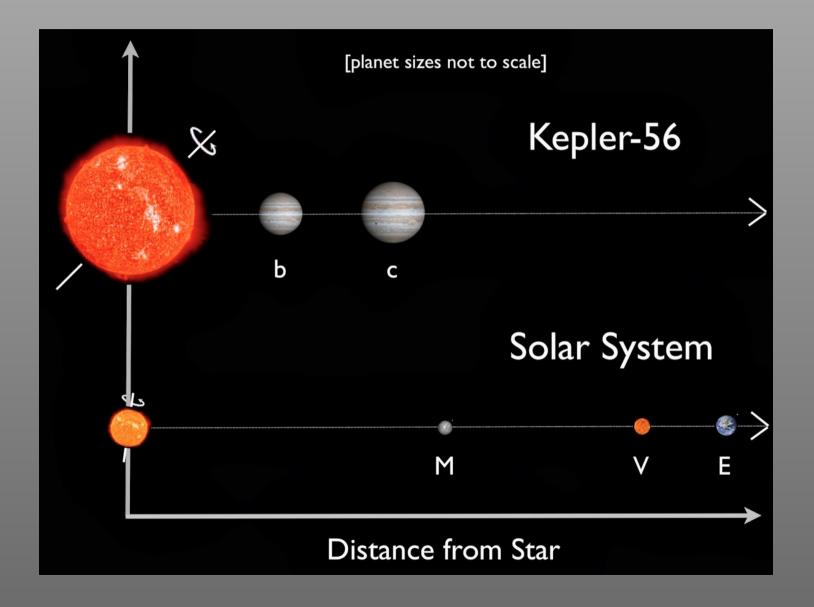
- growth limited to isolation mass
- planetesimals prone to fragmentation
- all growth is local
- planetesimal accr. slow in outer disk



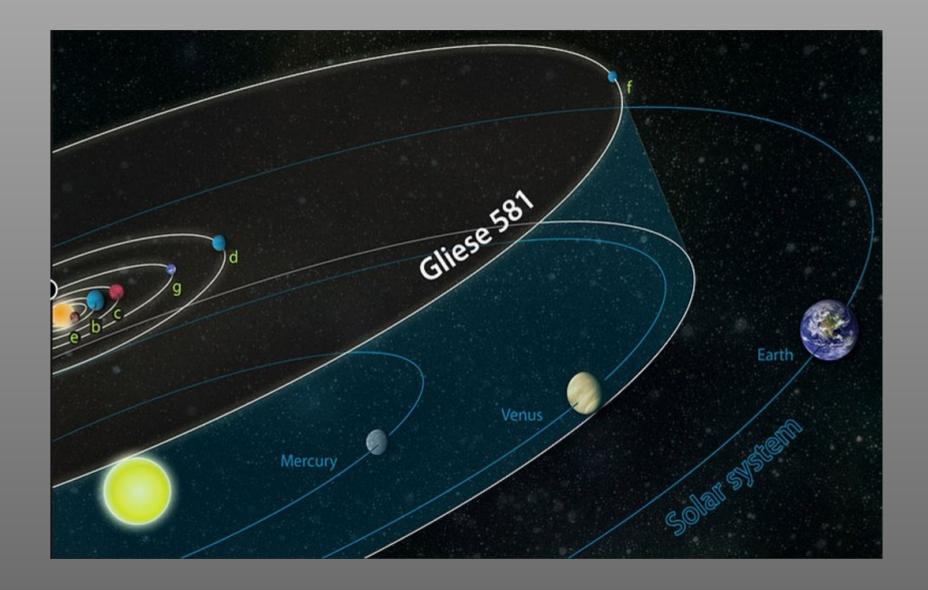
# Architecture solar system



# Kepler-56



# Gliese 581



# Kepler 11

