

## The Vectors of HGA $\mathbb{R}_{d,0,1}$

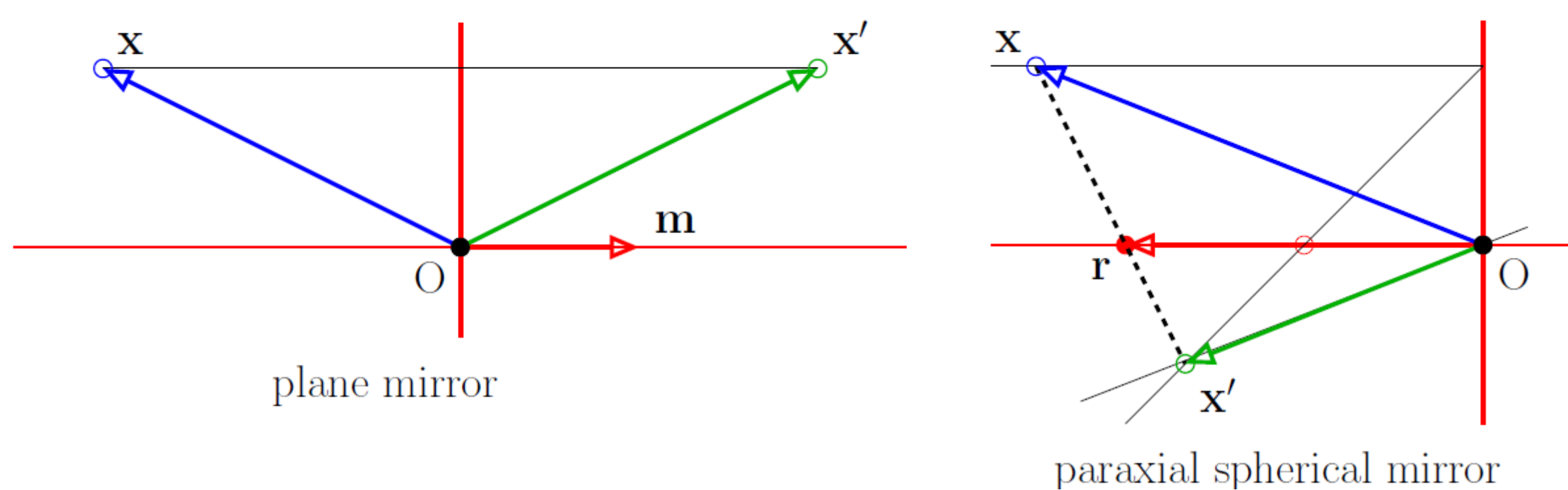
HGA  $\mathbb{R}_{d,0,1}$  is an algebra in which homogeneous points are represented as vectors. We choose an orthonormal basis  $\{e_0, e_1, \dots, e_d\}$  of anti-commuting ‘unit’ vectors:  $e_i e_j = -e_j e_i$  and  $e_0 e_i = -e_i e_0$ , with  $(e_0)^2 = 0$  and  $(e_i)^2 = 1$ .

Three kinds of vectors, differing in algebraic properties and geometric semantics:

- **Directions:** A purely Euclidean vector  $\mathbf{m}$  represents a 1-direction in  $d$ -space. Its inverse is  $\mathbf{m}^{-1} \equiv \mathbf{m}/(\mathbf{m} \cdot \mathbf{m})$ . Directions are ‘ideal points’ at infinity.
- **The origin:** The point at the origin is represented by the null vector  $e_0$ . Since  $e_0^2 = 0$ , it is *not* invertible, and therefore  $e_0$  cannot be used as a versor.
- **Non-origin points:** A general point at  $\mathbf{p}$  is of the form  $P = \alpha(e_0 + \mathbf{p})$ . Its inverse is  $P^{-1} = P/(P \cdot P) = P/(\alpha^2 \mathbf{p}^2)$ , iff  $\alpha \neq 0$  and  $\mathbf{p}^2 \neq 0$ . For versors we can use normalized points with  $\alpha = 1$ .

## The Vector Versors of HGA

$$-\mathbf{m}(e_0 + \mathbf{x})\mathbf{m}^{-1} = e_0 - \mathbf{m}\mathbf{x}\mathbf{m}^{-1} \quad -(\mathbf{e}_0 + \mathbf{r})(e_0 + \mathbf{x})(\mathbf{e}_0 + \mathbf{r})/\mathbf{r}^2 = (1 - 2\mathbf{r}^{-1} \cdot \mathbf{x})(e_0 + \frac{-\mathbf{r}\mathbf{x}\mathbf{r}^{-1}}{1 - 2\mathbf{r}^{-1} \cdot \mathbf{x}})$$



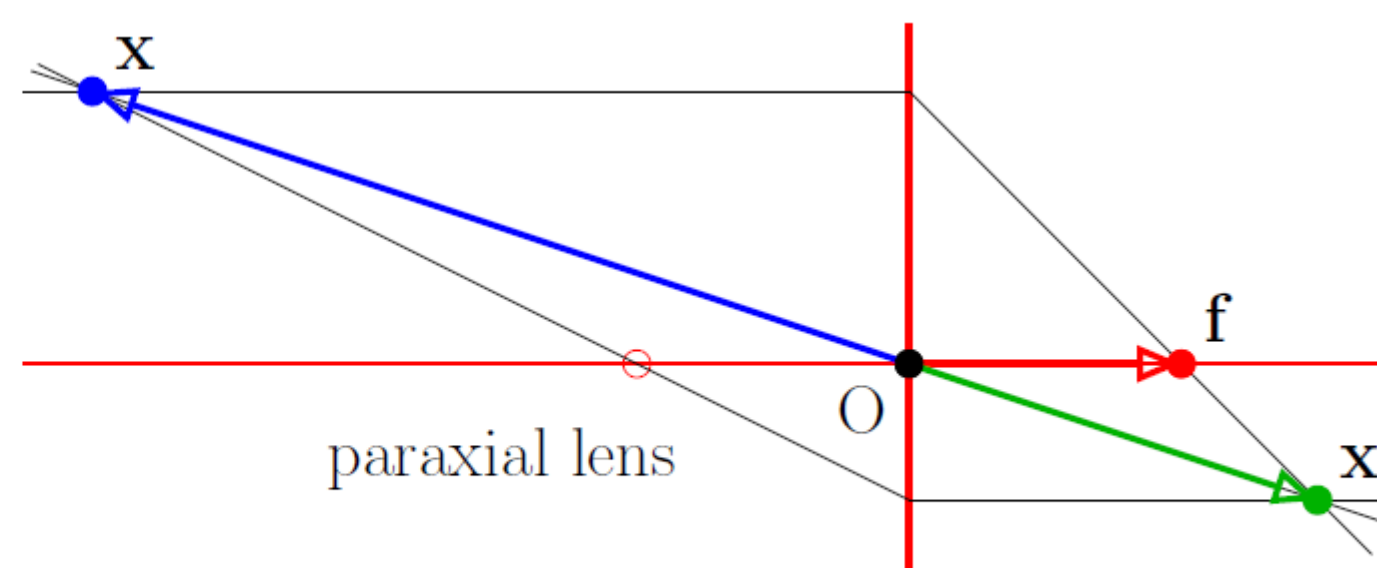
## The Lensing Versor $L_f$ in HGA

View a lens as the combination of a spherical mirror and a reflection, so use HGA versor  $\mathbf{m}R$  to represent it. We should take  $\mathbf{m} = \mathbf{r}$  and relate  $R$  to the focal point  $F$  by  $\mathbf{f} = -\mathbf{r}/2$ . That gives:

$$\mathbf{r}(e_0 + \mathbf{r}) = \mathbf{r}^2 - e_0 \mathbf{r} = \mathbf{r}^2(1 - e_0 \mathbf{r}^{-1}) \propto 1 - e_0 \mathbf{r}^{-1} = 1 + \frac{1}{2} e_0 \mathbf{f}^{-1} \equiv L_f.$$

Now sandwiching a point  $X = e_0 + \mathbf{x}$  yields:

$$\begin{aligned} L_f(e_0 + \mathbf{x})L_f^{-1} &= (1 + \frac{1}{2}e_0\mathbf{f}^{-1})(e_0 + \mathbf{x})(1 - \frac{1}{2}e_0\mathbf{f}^{-1}) \\ &= e_0 + \mathbf{x} + \frac{1}{2}e_0\mathbf{f}^{-1}\mathbf{x} - \frac{1}{2}\mathbf{x}e_0\mathbf{f}^{-1} \\ &= (1 + \mathbf{f}^{-1} \cdot \mathbf{x})e_0 + \mathbf{x} \\ &= (1 + \mathbf{f}^{-1} \cdot \mathbf{x})(e_0 + \frac{\mathbf{x}}{1 + \mathbf{f}^{-1} \cdot \mathbf{x}}). \end{aligned}$$

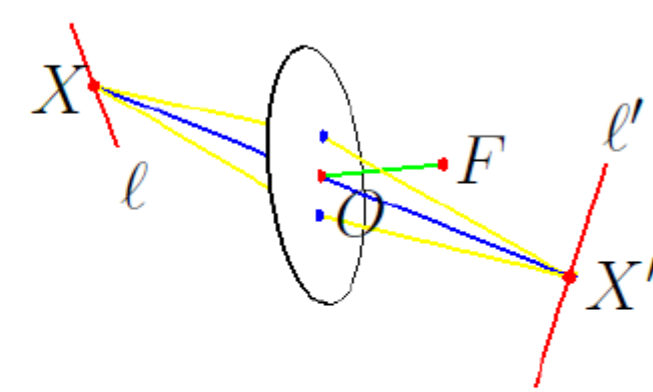


This is indeed the GA form of the familiar Euclidean lensing formula.

## HGA Bonus: Equivariance of Versor Mapping

The universality of the versor form allows direct transformation of lines and planes, e.g.

$$(L_f X_1 / L_f) \wedge (L_f X_2 / L_f) = L_f (X_1 \wedge X_2) / L_f.$$



However, a ray tracing matrix depends on the type of flat:

$$\text{point: } \begin{bmatrix} \mathbf{p} \\ 1 \end{bmatrix} \xrightarrow{L_f} \begin{bmatrix} [1] \\ [\mathbf{f}^{-1}]^T & 0 \\ 1 \end{bmatrix} \begin{bmatrix} \mathbf{p} \\ 1 \end{bmatrix} \quad \text{line: } \begin{bmatrix} \mathbf{u} \\ \mathbf{p} \times \mathbf{u} \end{bmatrix} \xrightarrow{L_f} \begin{bmatrix} [1] \\ [\mathbf{f}^{-1}]^T & 0 \\ 1 \end{bmatrix} \begin{bmatrix} \mathbf{u} \\ \mathbf{p} \times \mathbf{u} \end{bmatrix} \quad \text{plane: } \begin{bmatrix} \mathbf{n} \\ -\delta \end{bmatrix} \xrightarrow{L_f} \begin{bmatrix} [1] & -[\mathbf{f}^{-1}] \\ 0^T & 1 \end{bmatrix} \begin{bmatrix} \mathbf{n} \\ -\delta \end{bmatrix}.$$

Explicitly computing the matrices is better left to an HGA-to-LA compiler.

## Optics Anywhere, via Conformal GA

HGA null vector  $e_0$  for point at optical center is essential for lens versor. But there are no other null points in HGA  $\mathbb{R}_{d,0,1}$ , so we cannot translate by versors.

CGA  $\mathbb{R}_{d+1,1}$  (algebra of spheres) has all points as null vectors:  $x = n_o + \mathbf{x} + \frac{1}{2}\mathbf{x}^2 n_\infty$ . It also has translation versors. So let us embed HGA into CGA at each location!

## Changing Center by ‘from $c$ ’ Embedding $X|_c$

Choose any CGA point  $c$  as corresponding to the optical center  $e_0$  of HGA. An HGA point  $X$  at location  $\mathbf{x}$  viewed from this ‘origin’  $c$  is embedded as:

$$\text{‘from } c \text{’ map: } X|_c \equiv n_o + \mathbf{x} = n_o \cdot (-n_\infty \wedge \mathbf{x})$$

with  $c$  and  $x$  CGA points,  $n_o$  the CGA origin. You may pronounce ‘ $|_c$ ’ as ‘from  $c$ ’.

$$\text{structure-preserving: } x|_c \wedge y|_c = (x \wedge y)|_c.$$

Arbitrary HGA element  $X$  (point, line, plane, direction element) is represented as viewed ‘from  $c$ ’ by:

$$X \mapsto X|_c = c \cdot (-n_\infty \wedge X)$$

The ‘from  $c$ ’ mapping is ‘neopotent’:

$$\text{last application counts: } (X|_{c_1})|_{c_2} = X|_{c_2}.$$

We can therefore always re-represent a re-represented element. In a concatenation of lenses, there is *no need to revert* from  $X|_c$  to HGA  $X$  before the next step.

Jumping to a new optical center  $c$  is not a relative translation, but an absolute teleportation.

## The Lens Versor at Another Optical Center

HGA lens versor  $\exp(e_0 \wedge \mathbf{f}^{-1}/2)$  at origin  $e_0$ , move to  $C$ ? No HGA versor! CGA lens versor  $\exp(n_o \wedge \mathbf{f}^{-1}/2)$  at location  $n_o$ , move to  $c$ ? Have CGA versor!

CGA tangent vector  $n_o \wedge \mathbf{f}^{-1}$  at  $n_o$  moves to:  $c \wedge (\mathbf{f}^{-1} + (\mathbf{f}^{-1} \cdot \mathbf{c}) n_\infty)$ . With some algebra, we can rewrite this translated tangent multiplicatively:

$$T_c(n_o \wedge \mathbf{f}^{-1})T_c^{-1} = c/(c \wedge n_\infty \wedge \mathbf{f}) = c/(c \wedge n_\infty \wedge \mathbf{f}),$$

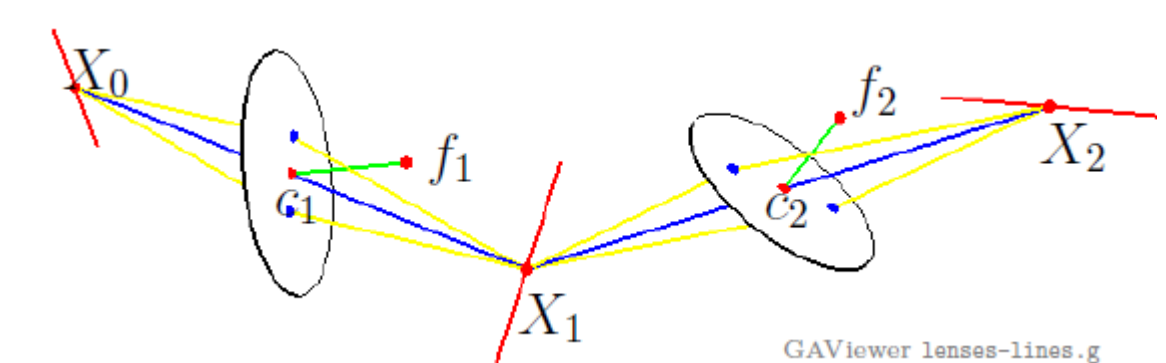
$f$  focal point, and  $\mathbf{f}$  relative vector from  $c$  to  $f$ . Indifferent to  $F$ -specification!

With that CGA tangent vector, the CGA lens versor with optical center  $c$  is:

$$\text{versor } L_{c,f} \text{ of lens at } c \text{ with focal point } f: L_{c,f} = \exp(c/(c \wedge n_\infty \wedge \mathbf{f})/2)$$

Apply the ‘at  $c$ ’ versor  $L_{c,f}$  to elements in the ‘from  $c$ ’ representation.

## Concatenation of Elements



1. Lenses with centers at points  $c_i$ , focal points at  $f_i$  (or use  $\mathbf{f}_i$ ). Form lens versors:

$$L_i \equiv \exp(\frac{1}{2} c_i / (c_i \wedge n_\infty \wedge \mathbf{f}_i)) = 1 + \frac{1}{2} c_i / (c_i \wedge n_\infty \wedge \mathbf{f}_i)$$

or the spherical mirror  $R_i = c_i - 2(c_i \wedge n_\infty) \cdot (c_i \wedge n_\infty \wedge \mathbf{f}_i)$ , or the planar mirror  $M_i = (c_i \wedge n_\infty) \cdot (c_i \wedge n_\infty \wedge \mathbf{f}_i)$ .

2. Embed the HGA element  $X$  into CGA by replacing its  $e_0$  by  $n_o$ . Then iterate:

$$X_0 = X, \quad X_i = \underline{L}_i[X_{i-1}|_{c_i}] = L_i(c_i \cdot (-n_\infty \wedge X_{i-1}))L_i^{-1} \quad \text{for } i = 1, \dots, n$$

or similarly for  $\underline{R}_i[]$  and  $\underline{M}_i[]$ , remembering to include the grade involution.

3. Final result may be expressed as  $X' = n_o \cdot (-n_\infty \wedge X_n)$  relative to an origin point  $n_o$ . It can be converted back to HGA by replacing  $n_o$  by  $e_0$ .

## Generating Ray Transfer System Matrices

- The HGA/CGA lens versor specification lenses any flat geometric primitive.
- The ‘from  $c$ ’ re-representation prepares it for the next optical element.
- Any flat element can be propagated through the optical system, so easy to find the total homogeneous transfer matrix for a composition of optical elements:
- Simply process appropriate  $i$ -th basis element by the iterative algorithm and denote the resulting components as  $i$ -th column of the transformation matrix. (E.g., matrix for imaging of arbitrary 3D line uses Plücker coordinate basis  $\{e_{01}, e_{02}, e_{03}, e_{23}, e_{31}, e_{12}\}$ .)
- Specification is immediately geometrical with  $(c_i, f_i)$  or  $(c_i, \mathbf{f}_i)$  pairs. Due to neopotent teleportation, no intermediate ‘relative Euclidean transformations’.
- This extends the 2D homogeneous matrix techniques of Corcovilos [2] to  $n$ -D.

Full paper [1] at:



l.dorst@uva.nl  
dorst.037@gmail.com  
(after 2024)

1. L. Dorst, “Paraxial geometric optics in 3D through point-based geometric algebra”, in CGI 2023, LNCS 14498, B. Sheng, J. Kim, N. Magnenat-Thalmann, and D. Thalmann, Eds. Springer Verlag, 2024, pp. 340-354. [https://dl.acm.org/doi/abs/10.1007/978-3-031-50078-7\\_27](https://dl.acm.org/doi/abs/10.1007/978-3-031-50078-7_27)
2. T. Corcovilos, “Beyond the ABCDs: A better matrix method for geometric optics by using homogeneous coordinates”, American Journal of Physics, vol. 91, pp. 449(457), 2023. <https://pubs.aip.org/aapt/ajp/article/91/6/449/2891435/>
3. L. Dorst, D. Fontijne, and S. Mann, Geometric Algebra for Computer Science. Morgan Kaufman, 2009. <https://www.geometricalgebra.org>
4. E. Lengyel, 2022. [Online]. <https://projectivegeometricalgebra.org/>
5. C. Gunn, “On the homogeneous model of Euclidean geometry”, in Guide to Geometric in Practice, L. Dorst and J. Lasenby, Eds. Springer-Verlag, 2011, pp. 297(327). [Online] <https://www.springer.com/gp/book/9780857298102>
6. C. Gunn, “Geometric algebras for Euclidean geometry”, Advances in Applied Clifford Algebras, vol. 27, pp. 185-208, 2017. <https://link.springer.com/article/10.1007/s00006-016-0647-0>