# Geometric Algebra and symmetry in crystallography and physics 

E. Hitzer ${ }^{\text {a }}$<br>${ }^{a}$ International Christian University<br>University of Amsterdam<br>Mitaka, Japan<br>hitzer@icu.ac.jp

## Summary of the Abstract

From viewpoints of crystallography and of elementary particles, we explore symmetries of multivectors in the geometric (Minkowski) algebra Cl(3,1) that can be used to describe space-time.

## Introduction

Recently, 9] and [8] classified multivectors based on their symmetries ${ }^{11}$ under space inversion (main grade involution in geometric algebra) $\hat{1}$, time reversal $1^{\prime}$ and reversion (called wedge reversion by them) $\tilde{1}$. [9] notes in the conclusions that One could perhaps explore charge reversal $(\hat{C})$, parity reversal $(\hat{P})$ and time reversal $(\hat{T})$ in the relativistic context [11]. In the standard model of elementary particle physics many experiments have confirmed violation of parity symmetry by the weak interaction, and of $\hat{C} \hat{P}$ symmetry. However, strong interactions by themselves do preserve $\hat{C} \hat{P}$ symmetry [5].
Therefore one aim of this research is to work in this direction by looking at the effect of these three symmetries on multivectors of $C l(3,1)$, a (geometric) algebra that can be used to express space-time physics, with $\hat{C}, \hat{P}$ and $\hat{T}$ symmetries, e.g., defined by [6], but here we only focus on the effect of these transformations on the 16 basis blades that constitute the multivector basis of $C l(3,1)$, and ignore any functional dependence of coefficients in linear combinations that might express spinors or other physical quantities. Furthermore, [11] and [6] have a clear preference for the use of $C l(1,3)$, while in this work we prefer ${ }^{2}$ to use $C l(3,1)$ because its volume-time subalgebra $\left\{1, e_{0}, e_{123}, e_{0123}\right\}$ isomorphic to quaternions, where $e_{0}$ expresses the time direction, at the foundation of the theory of space-time Fourier transforms [17].
[9] and [8] work signature independent for all Clifford algebras of quadratic spaces. We work in an algebra of specific signature and want to take advantage of the principal reverse ${ }^{3}$ (see e.g. [17] (2.1.12)), which applied to $C l(3,1)$ acts like the conventional reverse, but in addition changes the sign of the time vector $e_{0} \rightarrow-e_{0}$. So we focus on the group of eight symmetries generated by grade involution $\hat{1}$, reversion $\tilde{1}$ and principal reverse ${ }^{4} 1^{\prime}$.

[^0]Table 1: Action of involutions of group (1) on all 16 basis elements (2) of $C l(3,1)$. Tp. $=$ type with scalar $S$, time vector $V_{0}$ multiple of $e_{0}$, space vector $V$, bivector $B_{0}$ with $e_{0}$ factor, space bivector $B$, trivector $T_{0}$ with $e_{0}$ factor, space trivector $T$ and pseudoscalar quadvector $Q$. Bas. = basis element, $e=$ even, $o=$ odd.

| Tp. | Bas. | $\hat{1}$ | 1 | $\overline{1}$ | $1^{\prime}$ | $\hat{1}^{\prime}$ | $\tilde{1}^{\prime}$ | $\overline{1}^{\prime}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $S$ | 1 | $e$ | $e$ | $e$ | $e$ | $e$ | $e$ | $e$ |
| $V_{0}$ | $e_{0}$ | $o$ | $e$ | $o$ | $o$ | $e$ | $o$ | $e$ |
| V | $e_{1}$ | $o$ | $e$ | $o$ | $e$ | $o$ | $e$ | $o$ |
|  | $e_{2}$ | o | $e$ | o | $e$ | $o$ | $e$ | $o$ |
|  | $e_{3}$ | $o$ | $e$ | $o$ | $e$ | $o$ | $e$ | $o$ |
| $B_{0}$ | $e_{01}$ | $e$ | $o$ | $o$ | $e$ | $e$ | $o$ | $o$ |
|  | $e_{02}$ | $e$ | $o$ | $o$ | $e$ | $e$ | $o$ | $o$ |
|  | $e_{03}$ | $e$ | $o$ | $o$ | $e$ | $e$ | $o$ | $o$ |
| B | $e_{23}$ | $e$ | o | $o$ | $o$ | $o$ | $e$ | $e$ |
|  | $e_{31}$ | $e$ | o | o | o | $o$ | $e$ | $e$ |
|  | $e_{12}$ | $e$ | $o$ | $o$ | $o$ | $o$ | $e$ | $e$ |
| $T_{0}$ | $e_{023}$ | 0 | $o$ | $e$ | $e$ | $o$ | $o$ | $e$ |
|  | $e_{031}$ | 0 | $o$ | $e$ | $e$ | $o$ | $o$ | $e$ |
|  | $e_{012}$ | $o$ | o | $e$ | $e$ | $o$ | $o$ | $e$ |
| $T$ | $e_{123}$ | 0 | $o$ | $e$ | $o$ | $e$ | $e$ | $o$ |
| $Q$ | $e_{0123}$ | $e$ | $e$ | $e$ | $o$ | $o$ | $o$ | $o$ |

Our main reference is [22].

## Classification of multivectors in $C l(3,1)$

This talk is structured as follows. First we study the symmetries of $C l(3,1)$ multivectors under space inversion $\hat{1} M$, reversion $\tilde{1} M$, Clifford conjugation $\overline{1} M$ and principal reverse $1^{\prime} M$. These involutions generate by composition the following Abelian group of involutions

$$
\begin{equation*}
G=\left\{1, \hat{1}, \tilde{1}, \overline{1}, 1^{\prime}, \hat{1}^{\prime}, \tilde{1}^{\prime}, \overline{1}^{\prime}\right\} . \tag{1}
\end{equation*}
$$

Their action on the blade basis of the geometric algebra $C l(3,1)$ is listed in Table 1 . In the table horizontal lines separate the eight principle types of multivectors. Combinations of them yield 43 further types, i.e. a total of 51 types of multivector symmetry. The unit blade basis of the geometric algebra $C l(3,1)$ is given by one scalar, four vectors, six bivectors, four trivectors and one pseudoscalar quadvector $I$,

$$
\begin{equation*}
\left\{1, e_{0}, e_{1}, e_{2}, e_{3}, e_{01}, e_{02}, e_{03}, e_{23}, e_{31}, e_{12}, e_{023}, e_{031}, e_{012}, e_{123}, I=e_{0123}\right\} \tag{2}
\end{equation*}
$$

reversal of [9] and [8], which there also has the symbol $1^{\prime}$. Although we do obtain it for multivectors that have $e_{0}$ as a factor, by the product of reversion $\tilde{1}$ and principal reverse $1^{\prime}$ (our notation).

Table 2: Table of all compositions of symmetry operators $\hat{C}, \hat{P}$ and $\hat{T}$, where operations in the top row are applied first to $M$ followed by an operation from the first column. For example: combining $\hat{T} \hat{C}$ from the top row with $\hat{C} \hat{P}$ from the first column (6th row) shows that $\hat{C} \hat{P} \hat{T} \hat{C} M=\hat{P} \hat{T} M$.

| $\begin{aligned} & \text { 1st: } \\ & \text { 2nd: } \end{aligned}$ | 1 | $\hat{C}$ | $\hat{P}$ | $\hat{T} \hat{C}$ | $\hat{C} \hat{P}$ | $\hat{T}$ | $\hat{C} \hat{P} \hat{T}$ | $\hat{P} \hat{T}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 1 | $\hat{C}$ | $\hat{P}$ | $\hat{T} \hat{C}$ | $\hat{C} \hat{P}$ | $\hat{T}$ | $\hat{C} \hat{P} \hat{T}$ | $\hat{P} \hat{T}$ |
| $\hat{C}$ | $\hat{C}$ | 1 | $\hat{C} \hat{P}$ | - $\hat{T}$ | $\hat{P}$ | $-\hat{T} \hat{C}$ | $\hat{P} \hat{T}$ | $\hat{C} \hat{P} \hat{T}$ |
| $\hat{P}$ | $\hat{P}$ | $-\hat{C} \hat{P}$ | 1 | $\hat{C} \hat{P} \hat{T}$ | - $\hat{C}$ | $\hat{P} \hat{T}$ | $\hat{T} \hat{C}$ | $\hat{T}$ |
| $\hat{T} \hat{C}$ | $\hat{T} \hat{C}$ | $\hat{T}$ | $-\hat{C} \hat{P} \hat{T}$ | 1 | $\hat{P} \hat{T}$ | C | - $\hat{P}$ | $\hat{C} \hat{P}$ |
| $\hat{C} \hat{P}$ | $\hat{C} \hat{P}$ | $-\hat{P}$ | C | $\hat{P} \hat{T}$ | -1 | $\hat{C} \hat{P} \hat{T}$ | $-\hat{T}$ | $-\hat{T} \hat{C}$ |
| $\hat{T}$ | $\hat{T}$ | $\hat{T} \hat{C}$ | $\hat{P} \hat{T}$ | - $\hat{C}$ | $-\hat{C} \hat{P} \hat{T}$ | -1 | $\hat{C} \hat{P}$ | - $\hat{P}$ |
| $\hat{C} \hat{P} \hat{T}$ | $\hat{C} \hat{P} \hat{T}$ | $\hat{P} \hat{T}$ | $-\hat{T} \hat{C}$ | $\hat{P}$ | $\hat{T}$ | $-\hat{C} \hat{P}$ | -1 | - $\hat{C}$ |
| $\hat{P} \hat{T}$ | $\hat{P} \hat{T}$ | $\hat{C} \hat{P} \hat{T}$ | $\hat{T}$ | $\hat{C} \hat{P}$ | $-\hat{T} \hat{C}$ | $-\hat{P}$ | - $\hat{C}$ | -1 |

## Multivectors in $C l(3,1)$ under $\hat{C}, \hat{P}$ and $\hat{T}$ symmetries

Next, we consider aspects of charge conjugation, parity reversal and time reversal, when $C l(3,1)$

$$
\begin{equation*}
\hat{C} M=M e_{1} e_{0}, \quad \hat{P} M=e_{0} M e_{0}, \quad \hat{T} M=I e_{0} M e_{1} \tag{3}
\end{equation*}
$$

is applied in the description of elementary particle physics. The composition of the symmetries $\hat{C}, \hat{P}$ and $\hat{T}$ is shown in Table 2. The table does not change when the algebra $C l(1,3)$ is employed instead. Furthermore, the table is also the same, when the full $\hat{C}, \hat{P}$ and $\hat{T}$ symmetries are applied (including the operations on the spacetime vector argument of multivector functions as found in [6], equations (8.90)). An interesting feature is that Table 2 is isomorphic to the mutliplication table of the basis elements of $C l(3,0) \cong C l_{+}(3,1) \cong C l_{+}(1,3)$.
The application of the symmetries $\hat{C}, \hat{P}$ and $\hat{T}$ to the multivector basis of $C l(3,1)$ in Table 3. In contrast to Table 1, not only signs are changed, but also permutations occur limited to four subsets (including one subgroup) of four elements each.

## Summary

In this presentation we have pursued the application of elementary symmetries of the geometric algebra $C l(3,1)$ that can describe space-time. Inspired by [9] and [8], we chose three involutions of space inversion, reverse and principal reverse and studied the Abelian group thus generated and its action on the multivectors of $C l(3,1)$. We found that similar to [8], a classification in eight principal and further 43 types of multivectors is thus possible, leading to a total of 51 types. Then we looked at algebraic aspects of applying charge conjugation, parity reversal and time reversal to the multivector basis of $C l(3,1)$. We found that the composition of the symmetry operations $\hat{C}, \hat{P}$ and $\hat{T}$ forms an algebra isomorphic to $C l(3,0)$ and $C l_{+}(3,1)$, and we commented on the structures

Table 3: Application of charge conjugation $\hat{C}$, parity reversal $\hat{P}$ and time reversal $\hat{T}$ (top row) defined in (3), to all elements of the basis (first column) of $C l(3,1)$ given in (2).

| Basis | 1 | $\hat{C}$ | $\hat{P}$ | $\hat{T} \hat{C}$ | $\hat{C} \hat{P}$ | $\hat{T}$ | $\hat{C} \hat{P} \hat{T}$ | $\hat{P} \hat{T}$ |
| :--- | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 1 | $-e_{01}$ | -1 | $e_{0123}$ | $e_{01}$ | $e_{23}$ | $e_{0123}$ | $-e_{23}$ |
| $e_{0}$ | $e_{0}$ | $e_{1}$ | $-e_{0}$ | $e_{123}$ | $-e_{1}$ | $-e_{023}$ | $e_{123}$ | $e_{023}$ |
| $e_{1}$ | $e_{1}$ | $e_{0}$ | $e_{1}$ | $-e_{023}$ | $e_{0}$ | $e_{123}$ | $e_{023}$ | $e_{123}$ |
| $e_{2}$ | $e_{2}$ | $-e_{012}$ | $e_{2}$ | $-e_{031}$ | $-e_{012}$ | $e_{3}$ | $e_{031}$ | $e_{3}$ |
| $e_{3}$ | $e_{3}$ | $e_{031}$ | $e_{3}$ | $-e_{012}$ | $e_{031}$ | $-e_{2}$ | $e_{012}$ | $-e_{2}$ |
| $e_{01}$ | $e_{01}$ | -1 | $e_{01}$ | $-e_{23}$ | -1 | $-e_{0123}$ | $e_{23}$ | $-e_{0123}$ |
| $e_{02}$ | $e_{02}$ | $e_{12}$ | $e_{02}$ | $-e_{31}$ | $e_{12}$ | $-e_{03}$ | $e_{31}$ | $e_{03}$ |
| $e_{03}$ | $e_{03}$ | $-e_{31}$ | $e_{03}$ | $-e_{12}$ | $-e_{31}$ | $e_{02}$ | $e_{12}$ | $-e_{02}$ |
| $e_{23}$ | $e_{23}$ | $-e_{0123}$ | $-e_{23}$ | $-e_{01}$ | $e_{0123}$ | -1 | $-e_{01}$ | 1 |
| $e_{31}$ | $e_{31}$ | $-e_{03}$ | $-e_{31}$ | $-e_{02}$ | $e_{03}$ | $e_{12}$ | $-e_{02}$ | $-e_{12}$ |
| $e_{12}$ | $e_{12}$ | $e_{02}$ | $-e_{12}$ | $-e_{03}$ | $-e_{02}$ | $-e_{31}$ | $-e_{03}$ | $e_{31}$ |
| $e_{023}$ | $e_{023}$ | $e_{123}$ | $-e_{023}$ | $-e_{1}$ | $-e_{123}$ | $e_{0}$ | $-e_{1}$ | $-e_{0}$ |
| $e_{031}$ | $e_{031}$ | $e_{3}$ | $-e_{031}$ | $-e_{2}$ | $-e_{3}$ | $-e_{012}$ | $-e_{2}$ | $e_{012}$ |
| $e_{012}$ | $e_{012}$ | $-e_{2}$ | $-e_{012}$ | $-e_{3}$ | $e_{2}$ | $e_{031}$ | $-e_{3}$ | $-e_{031}$ |
| $e_{123}$ | $e_{123}$ | $e_{023}$ | $e_{123}$ | $e_{0}$ | $e_{023}$ | $-e_{1}$ | $-e_{0}$ | $-e_{1}$ |
| $e_{0123}$ | $e_{0123}$ | $-e_{23}$ | $e_{0123}$ | 1 | $-e_{23}$ | $e_{01}$ | -1 | $e_{01}$ |

found when $\hat{C}, \hat{P}$ and $\hat{T}$ are applied to the complete set of basis blades of $C l(3,1)$. Analogous results hold when the algebra $C l(1,3)$ is employed and when the full $\hat{C}, \hat{P}$ and $\hat{T}$ symmetries of [6], equations (8.90), are used. It may be interesting to apply both approaches in Clifford space gravity [3], and the study of elementary particles using a new embedding of octonions in geometric algebra [23, 21].

## Acknowledgments

The author wishes to thank God: Do not conform to the pattern of this world, but be transformed by the renewing of your mind. Then you will be able to test and approve what God's will is - his good, pleasing and perfect will. (Paul's recommendation in Romans 12:2, NIV). He further thanks his colleagues C. Perwass, D. Proserpio, S. Sangwine, the organizers of the AGACSE 2024 conference in Amsterdam, The Netherlands, and the anonymous reviewers of this abstract.

## References

[1] R. Abłamowicz, B. Fauser, On the transposition anti-involution in real Clifford algebras I: the transposition map, Linear and Multilinear Algebra, 59:12, pp. 13311358 (2011), DOI: $10.1080 / 03081087.2010 .517201$.
[2] S. Breuils, K. Tachibana, E. Hitzer, New Applications of Clifford's Geometric Algebra. Adv. Appl. Clifford Algebras 32, 17 (2022). DOI: https://doi.org/10.1007/ s00006-021-01196-7.
[3] C. Castro, Progress in Clifford Space Gravity, Adv. Appl. Clifford Algebras 23, pp. 39-62 (2013). DOI: https://doi.org/10.1007/s00006-012-0370-4.
[4] W. K. Clifford, Applications of Grassmann's Extensive Algebra, American Journal of Mathematics, Vol. 1, No. 4, pp. 350-358 (1978), DOI: https://doi.org/10. 2307/2369379.
[5] J. Dingfelder,T. Mannel, Mischung mit System, Physik Journal Vol. 22, No. 10, pp. 32-38 (2023)
[6] C. Doran, A. Lasenby, Geometric Algebra for Physicists, Cambridge University Press, Cambridge (UK), 2003.
[7] L. Dorst, D. Fontijne, S. Mann, Geometric algebra for computer science, an objectoriented approach to geometry, Morgan Kaufmann, Burlington (2007).
[8] P. Fabrykiewicz, A note on the wedge reversion antisymmetry operation and 51 types of physical quantities in arbitrary dimensions, Acta Cryst. A79, 381-384 (2023), DOI: https://doi.org/10.1107/S2053273323003303.
[9] V. Gopalan, Wedge reversion antisymmetry and 41 types of physical quantities in arbitrary dimensions, Acta Cryst. A76, pp. 318-327 (2020), DOI: https://doi. org/10.1107/S205327332000217X.
[10] D. Hestenes, J. W. Holt, Crystallographic space groups in geometric algebra, J. Math. Phys. 48(2) : 023514 (2007), DOI: https://doi.org/10.1063/1.2426416.
[11] D. Hestenes, Space-Time Algebra, Birkhäuser, Basel, 2015.
[12] D. Hildenbrand, Foundations of Geometric Algebra Computing, Springer, Berlin, 2013. Introduction to Geometric Algebra Computing, CRC Press, Taylor \& Francis Group, Boca Raton, 2019.
[13] E. Hitzer, C. Perwass, Interactive 3D Space Group Visualization with CLUCalc and the Clifford Geometric Algebra Description of Space Groups, Adv. Appl. Clifford Algebras 20, pp. 631-658 (2010). DOI: https://doi.org/10.1007/ s00006-010-0214-z.
[14] E. Hitzer, T. Nitta, Y. Kuroe, Applications of Clifford's Geometric Algebra. Adv. Appl. Clifford Algebras 23, 377-404 (2013). DOI: https://doi.org/10.1007/ s00006-013-0378-4
[15] E. Hitzer, Creative Peace License. http://gaupdate.wordpress.com/2011/12/ 14/the-creative-peace-license-14-dec-2011/, last accessed: 12 June 2020.
[16] E. Hitzer, Introduction to Clifford's Geometric Algebra. SICE Journal of Control, Measurement, and System Integration, Vol. 51, No. 4, pp. 338-350, April 2012, (April 2012). Preprint DOI: https://doi.org/10.48550/arXiv.1306.1660
[17] E. Hitzer, Quaternion and Clifford Fourier Transforms, Taylor and Francis, London, 2021.
[18] E. Hitzer, Book Review of An Introduction to Clifford Algebras and Spinors. By Jayme Vaz Jr and Roldao da Rocha Jr. Oxford University Press, 2019. Acta Cryst. A76, Part 2, pp. 269-272 (2020), DOI: https://doi.org/10.1107/ S2053273319017030.
[19] E. Hitzer, C. Perwass, Space Group Visualizer, Independently published - KDP, Seattle (US), 2021.
[20] E. Hitzer, C. Lavor, D. Hildenbrand, Current Survey of Clifford Geometric Algebra Applications. Math. Meth. Appl. Sci. pp. 1331-1361 (2022). DOI: https: //onlinelibrary.wiley.com/doi/10.1002/mma.8316.
[21] E. Hitzer, Extending Lasenby's embedding of octonions in space-time algebra $C l(1,3)$, to all three- and four dimensional Clifford geometric algebras $C l(p, q)$, $n=p+q=3,4$, Math. Meth. Appl. Sci. 47, pp. 140-1424 (2024), DOI: https://doi.org/10.1002/mma.8577.
[22] E. Hitzer, On Symmetries of Geometric Algebra Cl $(3,1)$ for Space-Time, Adv. Appl. Clifford Algebras 34:30, 14 pages (2024). DOI: https://doi.org/10.1007/ s00006-024-01331-0, Preprint: https://vixra.org/abs/2401.0125.
[23] A. Lasenby, Some recent results for $S U(3)$ and octonions within the geometric algebra approach to the fundamental forces of nature, Math. Meth. Appl. Sci. 47, pp. 1471-1491 (2024), DOI: https://doi.org/10.1002/mma. 8934 .
[24] P. Lounesto, Cliff. Alg. and Spinors, 2nd ed., CUP, Cambridge, 2006.
[25] S. J. Sangwine, E. Hitzer, Clifford Multivector Toolbox (for MATLAB), Adv. Appl. Clifford Algebras 27(1), pp. 539-558 (2017). DOI: https://doi.org/10. 1007/s00006-016-0666-x, Preprint: http://repository.essex.ac.uk/16434/ 1/author_final.pdf.
[26] B. Schmeikal, Minimal Spin Gauge Theory. Adv. Appl. Clifford Algebras 11, pp. 63-80 (2001), DOI: https://doi.org/10.1007/BF03042039.
[27] G. Schubring (ed.), H. G. Grassmann (1809-1877): Visionary Mathematician, Scientist and Neohumanist Scholar, Kluwer, Dordrecht, 1996.


[^0]:    ${ }^{1}$ Note that [9] and [8] use for space inversion $\overline{1}$, and for (wedge) reversion $1^{\dagger}$.
    ${ }^{2}$ Another notable work using $C l(3,1)$ in elementary particle physics is, e.g. [26].
    ${ }^{3}$ The principal reverse is in geometric algebra the equivalent of matrix transposition, see [1].
    ${ }^{4}$ The reader should be aware that therefore in this work we do not use a priori the notion of time

