GRAPPA seminar

Supernova remnants in gamma rays and acceleration of Galactic cosmic rays

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University of Amsterdam

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Outline

Supernovae

- Types of supernovae
- Supernova remnants
- 2 Cosmic Rays
- $\bigcirc \gamma$ Rays
 - Break



Propagation

- Ground Based Detection
 - first ground based detections
 - air showers
 - shower detection methods
 - air shower reconstruction
 - some HESS results
 - Break

Why supernovae?

Energy scale

Supernovae can release $10^{53} ergs \triangleq 10^{46} J,$ one of the most energetic processes in the known universe

Rate

More energetic events (e.g. gamma-ray bursts) are too rare Proof

Only proven source so far ...



[Vink(2012)]

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Core collapse supernovae

• Final stage of the evolution of massive ($M\gtrsim 8M_{\odot})$ stars

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Hydrogen	11 Myrs	Н	He
Helium	2 Myrs	He	C,O
Carbon	2000 yrs	С	Ne,Mg
Neon	2.6 yrs	Н	He
Oxygen	0.7 yrs	He	C,O
Silicon	18 days	Si,S,Ar,Ca	Fe,Ni,Cr,Ti

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- Every next step produces less energy per unit mass of fuel
- $^{56}\mbox{Fe}$ most stable element, no energy to be gained from fusing it \rightarrow star builds up an iron core

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Core collapse

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 - Photodisintegration starts destroying nuclei, particularly

$$^{56}_{26}$$
Fe + $\gamma \rightarrow 13^4_2$ He + 4n and 4_2 He + $\gamma \rightarrow 2p^+ + 2n$ (1)

Endothermic reactions: take up energy otherwise used to generate the pressure necessary to support the core

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Endothermic reactions: take up energy otherwise used to generate the pressure necessary to support the core

Free electrons supporting the core (electron degeneracy) are captured by heavy nuclei and free protons from photodisintegration, e.g.:

$$p^+ + e^- \to n + \nu_e \tag{2}$$

Neutrinos escape, carrying away enormous amounts of energy

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End result: star loses so much energy, it can no longer support itself
 → core collapses under its own gravity

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- Causes even higher temperatures \rightarrow more photodisintegration \rightarrow shockwave loses most of its energy and stalls
- However, photodisintegration creates expanding neutrinosphere which impacts the accretion shock, causing it to expand again

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SD When $M_{WD} > M_{Chandrasekhar} \simeq 1.4 M_{\odot}$, carbon fusion ignites in the core Core is degenerate \rightarrow doesn't expand so doesn't cool \rightarrow runaway reaction \rightarrow explosion

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 - DD Detonation, deflagration or delayed detonation

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- **1** Ejecta dominated: total swept up mass less than mass of the ejecta
- Sedov-Taylor: swept up mass dominates, but negligible radiative losses
- Pressure driven: radiative losses are significant
- Merging: remnant mixes with the interstellar medium

Acceleration mechanisms

- Many different acceleration mechanisms for galactic particles.
- Acceleration generally assumed in or near the source.
- We discuss 6 plausible mechanisms.

1 - Cyclotron

• Zeeman splitting observed in sunspot. Determine magnetic field strengths.

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1 - Cyclotron

- Zeeman splitting observed in sunspot. Determine magnetic field strengths.
- Caused by moving protons / electrons.

$$\phi = \int B \cdot dA = B\pi R^2$$



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Cosmic Rays

1 - Cyclotron

$$-\frac{d\phi}{dt} = \oint E \cdot ds = U$$

• Generation / decay of magnetic fields cause electric fields.

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Cosmic Rays

1 - Cyclotron

$$-\frac{d\phi}{dt} = \oint E \cdot ds = U$$

• Generation / decay of magnetic fields cause electric fields.

• Electric fields cause accelerations of protons / electrons.

Energy gained after one orbit is equal to
$$eU$$
.
 $R = 10^7$ m $\frac{dB}{dt} = 0.2$ Tesla/day
 $eU = 0.73$ GeV

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1 - Cyclotron

The Cyclotron mechanism provides the right energies (particles up to 100 GeV).

But stable circular orbits require additional forces!

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2 - Sunspot Pairs

Sunspots often come in pairs with opposite magnetic polarity. They approach each other and then merge.

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Sunspots often come in pairs with opposite magnetic polarity. They approach each other and then merge.

Electric field $E \sim v \times B$ up to 10 V/m.

- Distance of 10⁷ m
- B = 0.2 T
- $v = 10^7 \text{ m/day}$

can give energies in GeV range for protons.

Cosmic Rays

2 - Sunspot Pairs

Conclusion:

Same energies as Cyclotron, but doesn't require additional forces!

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3 - Shock Acceleration

Shock Wave Acceleration Situation 1: particle collides with shock front head on.

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$$\Delta E = \frac{1}{2}m(v + (u_1 - u_2)^2 - \frac{1}{2}mv^2$$

Linearly: $\frac{\Delta E}{E} \sim 2\frac{u_1 - u_2}{v}$
Relativistically: $\frac{\Delta E}{E} \sim \frac{4}{3}\frac{u_1 - u_2}{c}$

Situation 2: particle within the shock front bouncing between inner and outer front.

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(In galactic nuclei, u_2 can even go up to 0.9 c)

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Energies up to \sim 100 TeV can be explained!

(Linear shock acceleration mechanisms are also called Fermi mechanism of first order.)

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Cosmic/gamma-ray particles can interact with magnetic clouds.

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If v and u parallel: $\Delta E \sim \frac{1}{2}m(2uv + u^2)$ If v and u anti-parallel: $\Delta E \sim \frac{1}{2}m(-2uv + u^2)$ On average: $\Delta E \sim mu^2$ $\frac{\Delta E}{E} \sim 2\frac{u^2}{v^2}$

$$\Delta E \sim m u^2$$

Quadratic in cloud velocity, hence Fermi mechanism "of second order".

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Remains correct even relativistically.

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Since cloud moves slowly, energy gain per collision is small. Acceleration takes a long time.

Magnetic clouds have higher gas density and therefore higher interaction probability.

$\Delta E \sim mu^2$

Still, particles lose energy between collisions by interacting with (inter)galactic gas. Minimum initial energy is needed for particles.

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Still, particles lose energy between collisions by interacting with (inter)galactic gas. Minimum initial energy is needed for particles.

This minimum energy could be provided by head on collision mechanism (Fermi mechanism of first order).

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Pulsars: Spinning, magnetized neutron stars \sim 20 km radius.

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$$T_{\it pulsar} \sim T_{\it star} rac{R^2_{\it pulsar}}{R^2_{\it star}}$$

If $T_{star} \sim 1$ month, and $R_{star} \sim 10^6$ km, $T_{pulsar} \sim 1$ ms.

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$$v_{\it pulsar} = rac{2\pi R_{\it pulsar}}{T_{\it pulsar}} \sim 4\cdot 10^6 \,\,{
m m/s}$$

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Assuming magnetic flux is conserved:

$$B_{pulsar} = B_{star} \frac{R_{pulsar}^2}{R_{star}^2}$$

if
$$B_{star} = 0.1$$
 T, $B_{pulsar} = 2.5 \cdot 10^8$ T

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Single charged particles can gain 1000 TeV per meter!!

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Average pulsar accelerates $\sim 10^9$ years worth of particles.

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6 - Binaries

Binaries consisting of a normal star and a pulsar are an acceleration candidate as well.


Energy protons gain, falling in from infinity, is $E = G \frac{m_p M_{pulsar}}{R_{pulsar}} \sim 70$ MeV.

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The magnetic field of the neturon star produces a Lorentz force $F \sim evB \sim eE$.

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The magnetic field of the neturon star produces a Lorentz force $F \sim evB \sim eE$.

The particles obtain an energy from this force $E = \int F \cdot ds = evB\Delta s$ up to $\sim 10^{19}$ eV.

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Assumed that jets of highly relativistic particles are accelerated near a black hole / nucleus of a galaxy and injected into its radiation field.

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Both protons and electrons can produce high energy γ -rays by inverse Compton scattering off accelerated electrons.

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Both protons and electrons can produce high energy γ -rays by inverse Compton scattering off accelerated electrons.

As a consequence of this, high-energy neutrinos are created in the decays of charged pions. Detections can presumably be made only if the jets are directed at us.



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- O Sunspot Pairs

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- Fermi mechanism of second order (colliding with gas clouds)
- Pulsars
- Accretion disks of binary systems

 $\gamma~{\rm Rays}$

Gamma Ray Production (Omer)

 γ Ray Production Mechanisms Review

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γ Rays from Supernova Remnants

- γ -ray astronomy is important for determining the CR content of astrophysical sources
- In γ -rays one can observe photons emitted as a result of hadronic CRs, which make up 99% of the CRs observed on Earth
- Three different particle radiation processes are considered most dominent in the supernova spectrum at γ-ray energy scales:
 - nuclear pion-production interactions
 - nonthermal electron bremsstrahlung
 - Compton scattering

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 $\gamma~{\rm Rays}$

Model of Suprenova Explosion

The supernova explosion is modeled as an expanding spherical shell of material that sweep up matter from the surrounding interstellar medium (ISM)



$$E_{ke} = \frac{1}{2} [M_0 + M_{su}(t)] v^2$$
(3)

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$$\frac{dE_{ke}^{su}(t)}{dt} = \frac{d}{dt} \left[\frac{1}{2}M_{su}v^2\right] = \frac{1}{2}\dot{M}_{su}v^2 + M_{su}v\dot{v}$$
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Density of the target particles for interactions

There are three targets available for particle-particle interactions

- Explosion mass : $n_{ex}(t) = M_0/m_p V_{sh}(t)$
- Swept-up mass : increased by the comprassion ratio, for strong shock the shell of swept-up mass has density $n_{su} \cong 4n_0$
- ISM gas *n*₀

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 $\gamma~{\rm Rays}$

Nuclear pion-production interactions

Rate of change of the Lorentz factor

$$\begin{aligned}
-\dot{\gamma}_{pp} &= K_p c \sigma_{pp} n(t) \gamma_p \\
\text{Energy-loss time scale} \\
\hline t_{pp} &= |\gamma_p/\dot{\gamma}_{pp}| = [K_p c \sigma_{pp} n(t)]^{-1} \cong 2.2 \times 10^{15}/n(t) \text{ s} \\
\text{Expected luminosity} \\
\hline L_{pp}(t) \cong \frac{\eta_{pp} E_{su}(t)}{3t_{pp}(t)}
\end{aligned}$$

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 $\gamma~{\rm Rays}$

Nonthermal electron bremsstrahlung

Rate of change of the Lorentz factor
$$\dot{k} = k \cos \alpha z \left[\sum_{n=2}^{\infty} \frac{7}{2} \left(\frac{7}{2} + 1 \right) \right]_{n=1}^{\infty}$$

$$\frac{-\gamma_{ff} = \kappa_{ff} \alpha_f c \sigma_{T} [\Sigma n_Z \Sigma (\Sigma + 1)] \gamma_e}{\text{Energy-loss time scale}}$$

$$t_{ff}(t) = |\gamma_e/\dot{\gamma}_{ff}| \cong 8.0 imes 10^{14}/n(t) \ s$$

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Compton scattering

Rate of change of the Lorentz factor

$$\begin{bmatrix}
-\dot{\gamma}_{C} = (4/3)c\sigma_{T}U_{\gamma}\gamma_{e}^{2}/m_{e}c^{2} \\
\text{Energy-loss time scale}
\end{bmatrix}$$

$$\begin{bmatrix}
t_{C}(s) = |\gamma_{e}/\dot{\gamma}_{C}| = 7.7 \times 10^{19}/\gamma_{e} s \\
\text{Expected luminosity} \\
\begin{bmatrix}
L_{C}(t) \cong \frac{\eta_{C} E_{su}(t)}{t_{C}(t)}
\end{bmatrix}$$

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 $\gamma~{\rm Rays}$

Energy-loss time scales and γ -ray luminosities



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SNR seen on Earth in γ -Rays

- Q_0^{π} , $Q_0^{
 m brem}$, and $Q_0^{
 m IC}$ gamma ray emissivities (cm $^{-3}$ s $^{-1}$ GeV $^{-1}$)
- Then the gamma ray flux observed at Earth, a distance *d* from the SNR, is given by -

$$F_{\gamma}(E_{\gamma},\alpha) = \frac{n_1 A_1 V}{4\pi d^2} \left[Q_0^{\pi}(E_{\gamma},\alpha) + R_e Q_0^{\text{brem}}(E_{\gamma},\alpha) + \frac{R_e}{n_1} Q_0^{\text{IC}}(E_{\gamma},\alpha) \right]$$
(5)

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Break

Coffee Break and Discussion

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Propagation Through Space

Cosmic RaysGamma Rays

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• standard non-linear diffusive shock acceleration (DSA) gives steepest spectrum $E^{-2} \rightarrow$ observed spectrum $E^{-2.7}$

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- need to account for interaction with galactic magnetic field turbulence
- $\bullet\,$ galactic magnetic field amplification $\rightarrow\,$ efficient CRs scatter back and forth the SNR shock
- self-consistent model:
 efficient CRs acceleration → magnetic field amplification
 feedback from amplified fields → efficient CR acceleration

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general procedure

- specify the CRs sources
- define the CRs halo shape and boundary conditions (CRs freely exit into intergalactic space)
- account for energy loss or gain processes in interstellar medium (ISM)
- nuclear fragmentation
- radioactive decay of unstable nuclei

• steady state transport equation in ISM

$$-\nabla D\nabla \Psi + \nabla (\mathbf{u}\Psi) - \frac{\partial}{\partial p} \left[p^2 K \frac{\partial}{\partial p} (p^{-2}\Psi) \right] \\ - \frac{\partial}{\partial p} \left(p \frac{\nabla \mathbf{u}}{3} \Psi \right) + \frac{\partial}{\partial p} (\dot{p}_{loss}\Psi) + \frac{\Psi}{\tau} = q \quad (6)$$

• diffusion equation

$$D = \frac{v r_g B^2}{12\pi k_{res} W(k_{res})} = \frac{v r_g^a}{3(1-a) k_L^{1-a}} \frac{B^2}{\delta B_L^2}$$
(7)

• GALPROP code for numerical simulation: solves (1) for nuclei, \bar{p}, e^-, e^+ ; computes γ -rays and synchrotron emission

• D determines the propagation of CRs in galactic magnetic fields

Image: A matrix

- D determines the propagation of CRs in galactic magnetic fields
- wave-particle interaction is of resonant character

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isotropization due to large scale ($\sim 100 pc)$ galactic magnetic field fluctuations

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isotropization due to large scale ($\sim 100 pc)$ galactic magnetic field fluctuations

- the spectrum of the MHD turbulence, $W(k_{res})$, determines the diffusion coefficient in (2)
- two proposed spectra

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MHD turbulence spectrum

- Kolmogorov spectrum $W(k) \propto k^{-\frac{5}{3}}$
 - accounts for reacceleration of CRs by MHD
 - leads to $D \propto v (p/Z)^{1/3}$
- Iroshnikov-Kraichnan spectrum $W(k) \propto k^{-rac{3}{2}}$
 - reacceleration with wave damping
 - leads to $D \propto v(p/Z)^{1/2}$

secondary-to-primary ratio

- $\bullet\,$ key information on CRs propagation from the abundance of light elements: $^{2}\mathrm{H},~^{3}\mathrm{H},$ Li, Be, B
- produced by spallation of heavier primary with ISM
- \bullet estimate secondary-to-primary ratio \rightarrow B/C
- allows to infer the MHD turbolence spectrum

Propagation

B/C ratio

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• $\rho_{CRs} \simeq \rho_{MHD}$

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 → integral part of CRs acceleration

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- Parker instability: short wavelenght MHD instability
- both lead to amplification of magnetic fields at the supernova shock
 → integral part of CRs acceleration
- Galactic gas halo might not be static (galactic wind) → CRs exiting the Galaxy increase the MHD turbulance → self-consistently determines diffusion-convection of CRs

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Photon propagation

- γ -rays propagate freely propagate through ISM
- their path might be deflected by gravitational lensing
- they "feel" the ISM magnetic field through the Faraday Effect

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Observations

Space-born detectors: Fermi- Large Area Telescope (LAT) Astro-rivelatore Gamma a Immagini Leggero (AGILE) AGILE and Fermi-LAT investigate complementary energy bands

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Tycho SNR

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- no leptonic model \rightarrow hadronic model (π^0 decay)

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Propagation

SNRs and Molecular clouds interaction



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- SNRs interacting with nearby Molecular Clouds (MCs) are the most luminous in γ -rays
- interaction evidenced by OH maser emission at 1720 MHz
- best observed SNRs: W44 and IC443
- in both γ-ray spectrum peaked at 1Gev lower energy cut-off below 200 MeV higher energy cut-off above 200 GeV
- \bullet fit the Fermi-LAT and AGILE data with π^0 decay spectrum

Propagation

W44 and IC443



first detections of SN

	length of				Historical Records		
date	visibility	remnant	Chinese	Japanese	Korean	Arabic	European
AD1604	12 months	$G4 \cdot 5 + 6 \cdot 8$	few	-	many	-	many
AD1572	18 months	$G120 \cdot 1 + 2 \cdot 1$	few	-	two	-	many
AD1181	6 months	3C58	few	few	-	-	-
AD1054	21 months	Crab Nebula	many	few	-	one	-
AD1006	3 years	SNR327.6+14.6	many	many	-	few	two
AD393	8 months	-	one	-	-	-	-
AD386?	3 months	-	one	-	-	-	-
AD369?	5 months	_	one	-	-	-	-
AD185	8 or 20 months	_	one	-	-	_	-

We will take another look at the remnants of the highlighted supernovas at the end of this presentation.

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the electromagnetic spectrum observed so far



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characterization of airshowers

- amount of particles
- direction of main axis
- spacial structure
- spread in time
- hadron content
- fluctuations in development
- muon content



air showers

different signatures

- Cherenkov light
- particles reaching the ground
- radio emission (as well from interaction with the geomagnetic field)
- air fluoresence
- acoustic effects



particle detector arrays

- sampling on arrival at ground level
- some information on the state of the shower from the arrival sequence of the current particle generation
- direction information from the charge separation of the magnetic field and the geometry / arrival times of the signal



example: Akeno Grand Air Shower Array in Japan

- takeing data since 1991
- 111 detectors
- approx 1km spacing
- measured gamma rays above the GZK-cutoff

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fly's eyes: air fluoresence detectors

- shower excites nitrogen in the atmosphere
- isotropic emission of fluoresence light (300-400 nm band)
- detection by PMTs
- advantage: able to monitor large areas and therefore aimed to detect rare ultra high gamma ray events



HiRes detector in Utah

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Cherenkov telescopes: review of Cherenkov light



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Cherenkov telescope: H.E.S.S.

goal: restruct energy, species and direction of the initial particle





Cosmic Ray Simulations for Kascade



The CORSICA code:

- interactions and cross-sections
- decay and propagation (ionization of the atmosphere, energy loss)
- (seasonal) composition of atmoshere, different layers
- earth magnetic field
- A word of caution:
 - knowledge of high energy interactions incomplete
 - extreme forward direction not accesible at colliders

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simulation results - photon, proton and iron at 10^{13} eV



CORSIKA simulation 100TeV gamma ray

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RCW86 possibly the remnant of SN AD185



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SN AD1006 remnant morphology



excess counts of gamma rays (with energy above 260 GeV) white region shows the earlier measured X-ray distribution

significance in standard deviations white region contains 80% of the respective X-ray energy

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SN AD1006 remnant gamma source



SHALL WE DO ONE SLIDE WITH ALL CONCLUSIONS FROM THE SECOND PART FOR THE DISCUSSION?

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Another Coffee? And Discussion..

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The End

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